Self adjoint operators and their spectra

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Our interest will be in self-adjoint operators A on Hilbert space \mathcal{H} . More specifically, we will be concerned with *bounded* self-adjoint A which admit a cyclic vector, i.e., a vector Ω such that the set $\{\Omega, A\Omega, \dots, A^n\Omega, \dots\}$ linearly spans a dense subspace of \mathcal{H} .

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It is a consequence of the Hermitian symmetry of any matrix representing A that Ω is a cyclic vector for A if and only if there exists an orthonormal basis $\{\Omega = e_0, e_1, e_2, \cdots\}$ for \mathcal{H} with respect to which the matrix representing A has the tridiagonal form

$$\begin{bmatrix} a_0 & b_0 & 0 & \cdots & \cdots \\ b_0 & a_1 & \overline{b_1} & \cdots & \cdots \\ 0 & b_1 & a_1 & \overline{b_2} & \cdots \\ \vdots & \vdots & \ddots & \ddots & \ddots \end{bmatrix}$$

with all the off-diagonal terms b_i being non-zero.

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Theorem

If μ is a probability measure on a compact set $\Sigma \subset \mathbb{R}$, the equation (Mf)(x) = xf(x) deines a bounded self-adjoint (multiplication) operator on $L^2(\Sigma, \mu)$, and the function $f_0 \equiv 1$ is a cyclic vector for M.

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A unit vector $\Omega \in \mathcal{H}$ is a cyclic vector for a bounded self-adjoint operator A on \mathcal{H} if and only if there exists a (uniquely determined) probability measure μ defined on the Borel subsets of $\Sigma = sp(A)$ and a unitary operator $U : \mathcal{H} \to L^2(\Sigma, \mu)$ such that $UAU^* = M$ and $U\Omega = f_0$.

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The measure μ above, is called the scalar spectral measure of A associated to $\Omega.$

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A compactly supported measure μ is determined - thanks to Weierstrass' polynomial approximation theorem and Riesz representation theorem - by the sequence $\{m_n : n = 0, 1, 2, \dots\}$ of **moments**, defined by

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We shall illustrate these notions and the way to use the one to get a description of the others by considereing one classic example.

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The standard semi-circular operator

The tridiagonal matrix

$$\left[\begin{array}{ccccccc} 0 & 1 & 0 & 0 & \cdots \\ 1 & 0 & 1 & 0 & \cdots \\ 0 & 1 & 0 & 1 & \cdots \\ \vdots & \vdots & \ddots & \ddots & \ddots \end{array}\right](*)$$

clearly represents a self-adjoint operator A_0 on ℓ^2 with cyclic vector e_0 .

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To calculate the moments $m_n = \langle A_0^n e_0, e_0 \rangle$, notice first that $A_0 = S + S^*$, where S (resp. S^*) denotes the right (resp., left-) shift on ℓ^2 with

$$Se_n = e_{n+1}, S^*e_{n+1} = e_n, S^*e_0 = 0$$

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It follows that $m_{2n+1} = 0$ and m_{2n} is the (Catalan) number C_n of (Dyck) paths of 2n steps from (0,0) to (2n,0), such that each step is from a (k, l) to a $(k, l \pm 1)$, and such that the path never goes below the x-axis. The Catalan numbers satisfy and are determined by the recurrence relation:

$$C_0 = 1$$

 $C_n = \sum_{p=1}^n C_{p-1} C_{n-p}$

The Cauchy transform of a measure μ supported on a compact subset of \mathbb{R} is the obviously holomorphic function defined on $\mathbb{C}^+ = \{z \in \mathbb{C} : Im \ z > 0\}$ by

$$G_{\mu}(z) = \int rac{d\mu(t)}{z-t}$$

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We then use the recursion relation satisfied by the Catalan numbers to compute the Cauchy transform of the scalar spectral measure of our semi-circular operator as follows:

Computation of G_{μ}

$$G_{\mu}(z) = \sum_{n=0}^{\infty} \frac{C_n}{z^{2n+1}}$$

$$= \frac{1}{z} + \sum_{n=1}^{\infty} \frac{1}{z^{2n+1}} C_n$$

$$= \frac{1}{z} + \sum_{n=1}^{\infty} \frac{1}{z^{2n+1}} \left(\sum_{p=1}^n C_{p-1} C_{n-p} \right)$$

$$= \frac{1}{z} + \frac{1}{z} \sum_{n=1}^{\infty} \sum_{p=1}^n \frac{C_{p-1}}{z^{2p-1}} \frac{C_{n-p}}{z^{2n-2p+1}}$$

$$= \frac{1}{z} + \frac{1}{z} \sum_{p=1}^{\infty} \frac{C_{p-1}}{z^{2p-1}} \left(\sum_{n=p}^{\infty} \frac{C_{n-p}}{z^{2n-2p+1}} \right)$$

$$= \frac{1}{z} (1 + G_{\mu}(z)^2)$$

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Hence $(G_{\mu}(z))^2 - zG_{\mu}(z) + 1 = 0$, and so

$$G_{\mu}(z)=\frac{z-\sqrt{z^2-4}}{2}$$

for a suitable branch of the complex square root, which must satisfy the obvious necessary condition that $\lim_{|z|\to\infty|} |G_{\mu}(z)| = 0$.

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It may be verified that the above limiting condition will be satisfied if the branch of $\sqrt{z^2 - 4}$ is defined as $\sqrt{z - 2} \times \sqrt{z + 2}$, where $\sqrt{z \mp 2} = \sqrt{|z \mp 2|} \exp(\frac{1}{2} \arg(z \mp 2))$, and this is defined only for $\arg(z \mp 2) \neq \frac{-\pi}{2}$,

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It may be further seen that this specification leads to a continuous determination of $G_{\mu}(z)$ for $Im \ z \ge 0$, in such a way that

$$G_{\mu}(t) = \left\{egin{array}{c} rac{t-\sqrt{t^2-4}}{2} & ext{if } t \geq 2 \ rac{t-i\sqrt{4-t^2}}{2} & ext{if } |t| \leq 2 \ rac{t+\sqrt{t^2-4}}{2} & ext{if } t \leq -2 \end{array}
ight.$$

The final step in the derivation of μ from G_{μ} is completed by the following application of the so-called *Stieltje's inversion formula* which shows that:

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From G_{μ} to μ

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In case $G_{\mu}(z)$ admits a continuous extension to $\mathbb{C}^+ \cup \mathbb{R}$ and if $g(t) = \lim_{\epsilon \to 0} G_{\mu}(t + i\epsilon)$, then

$$d\mu(t)=rac{-1}{\pi}(\textit{Im }g(t))dt$$

and we finally deduce that the scalar spectral measure of our semi-circular operator is the celebrated **Wigner semi-circular distribution** given by

$$\mu(E) = \frac{1}{2\pi} \int_{E \cap [-2,2]} \sqrt{4 - t^2} dt$$

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Lectures in the Combinatorics of Free Probability by Alexandru Nica and Roland Speicher, LMS Lecture Note Series 335, Cambridge University Press.

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By a not dissimilar analysis, we proved that this was a so-called *Free Poisson distribution*.

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