

# *Hysteresis in condensed matters: Three examples.*

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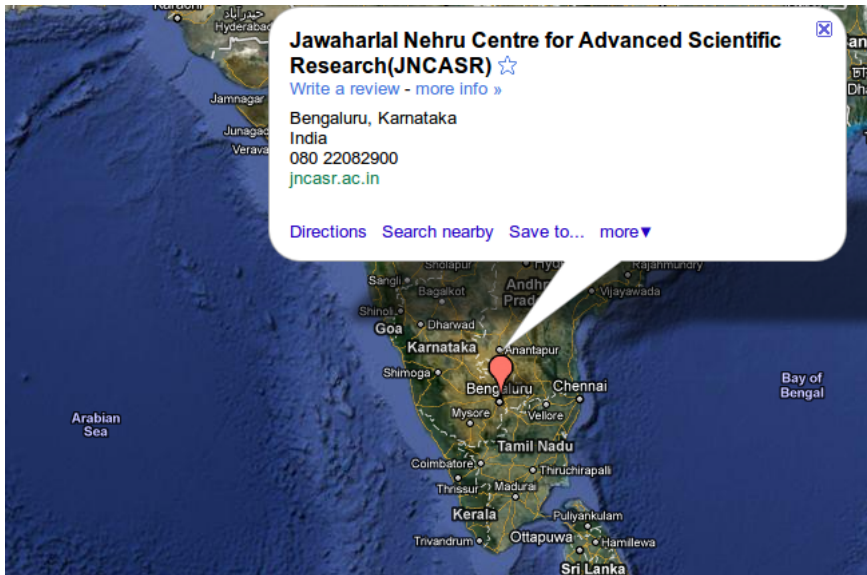
Condensed Matter Theory Seminar, Physics Dept., Purdue  
University

April 11, 2011





## My Place





## *People I have been working with*

**Ph.D. supervisor:** Dr. N. S. Vidhyadhiraja ([JNCASR](#))



**Collaborators:** Prof. Erica W. Carlson ([Purdue](#)), Prof. Karin A. Dahmen ([UIUC](#))





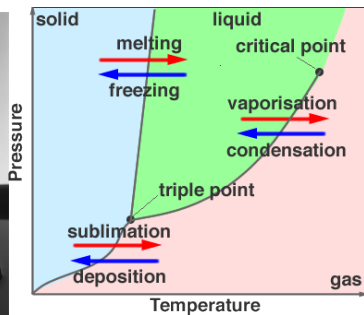
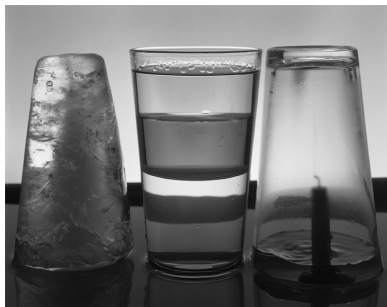
## *Quick outline*

1. Hysteresis in three phases: Liquid-gas-solid, Ferromagnets, Metal-Insulator (Mott)
2. Qualitative explanation
3. Further investigation: Avalanches



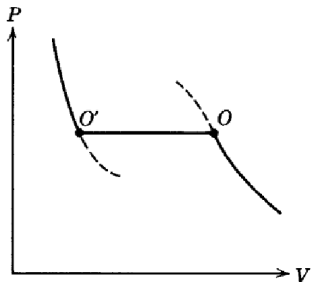
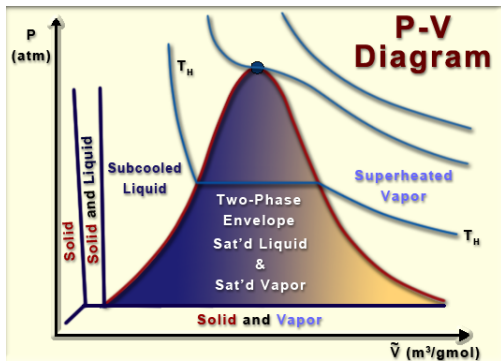
# *Liquid-gas-solid*

## Phase diagram



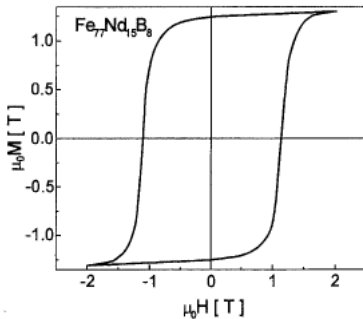
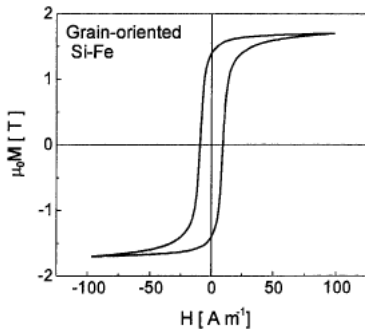


## Liquid-gas-solid





# Ferromagnets

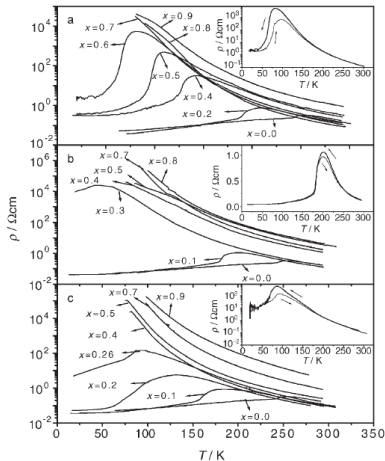


Ref. *Giorgio Bertotti, HYSTERESIS IN MAGNETISM*



# Metal-insulator transition

**Manganites:**  $(\text{La}_x\text{Ln}_{1-x})_{0.7}\text{Ca}_{0.3}\text{MnO}_3$



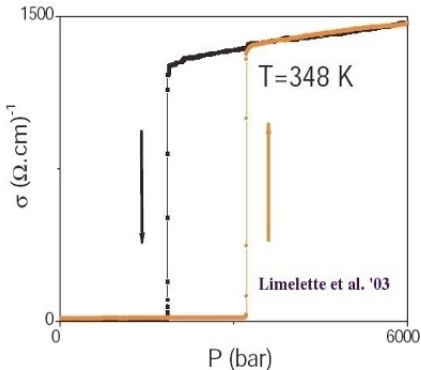
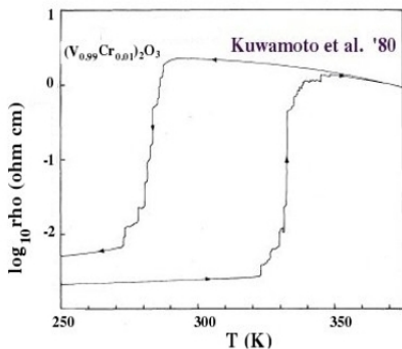
Shenoy, Sharma, Rao, *ChemPhysChem* 7, 2053 ('06)





## Metal-insulator transition

$V_2O_3$ : Contribution from Purdue University, *J. M. Honig's group*

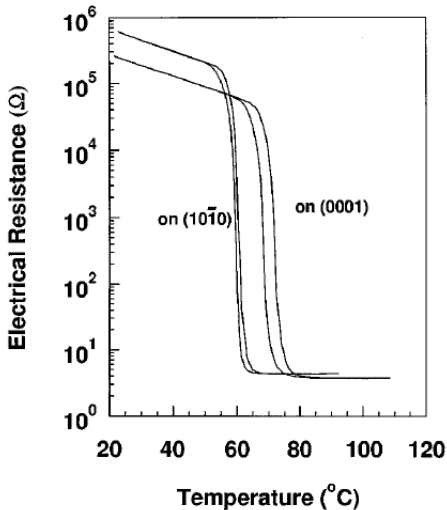


Temperature and pressure driven hysteresis



## *Metal-insulator transition*

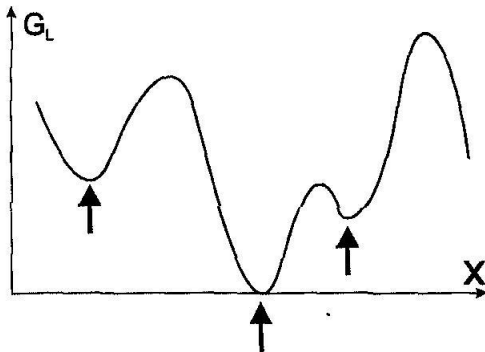
VO<sub>2</sub>: Kim and Kwok, APL ('94)





## *General theory for hysteresis*

Metastable states in free energy profile





## *Universality ?*

Can think of Ising-universality when we have two-state phases:  
Liquid and Gas  $\leftrightarrow$  Spin- $\uparrow$  FM and Spin- $\downarrow$  FM  $\leftrightarrow$  Metal and  
Insulator

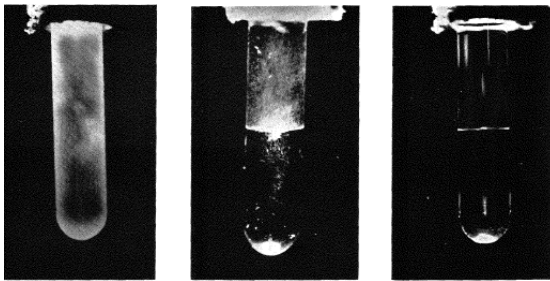
However, liquid-gas may differ since surface tension is an additional  
driving force.

Why thermal hysteresis is not common in ferromagnets?



## *Two-phase state: Evidence for coexistence*

### Liquid-gas



Ref. E. G. Stanley, *INTRO. TO PHASE TRANSITION AND CRITICAL PHENOMENA*



## *Two-phase state: Evidence for coexistence*

### Ferromagnet

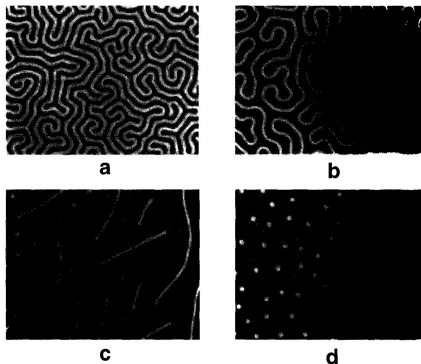
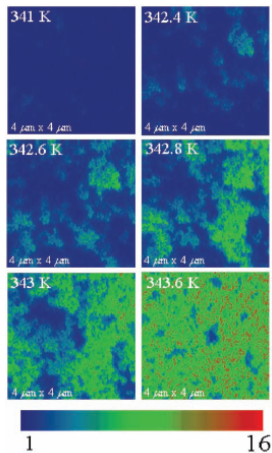


FIGURE 7.14. Domains in epitaxial garnet crystal  $6\text{ }\mu\text{m}$  thick, observed by Faraday effect. (a) Stripe domains under zero field. (b) Stripe domains under an external field perpendicular to the plane of the film. (c) and (d) Formation of bubble domains. Magnification is 300x (after [B.87], p. 166, Fig. 4.31, reprinted with kind permission from Academic Press).



## *Two-phase state: Evidence for coexistence*

**VO<sub>2</sub>:** Near-field Infra-red spectroscopy



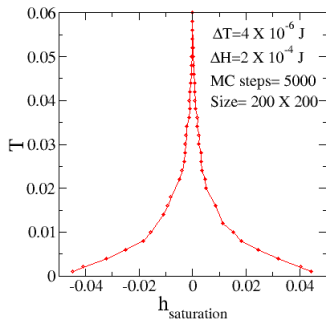
Ref. Qazilbash et al., *Science* 318, 1750 ('07)



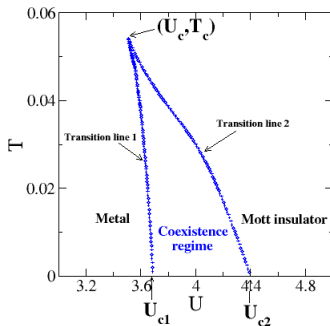
## Thermal hysteresis: Role of phase geometry

Ising model Vs Hubbard model (Dynamical mean-field th. scheme,  
Barman and Vidhyadhiraja, <http://arxiv.org/abs/1011.4478>)

Ising model phase diagram



HM (DMFT) phase diagram



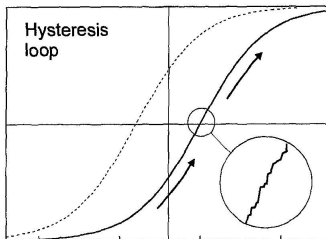
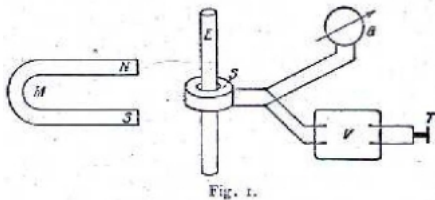
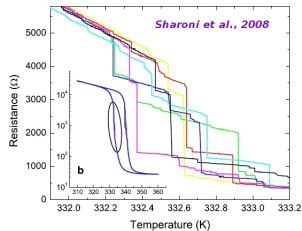
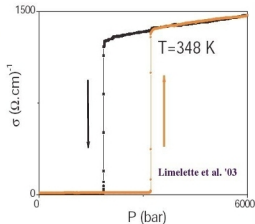
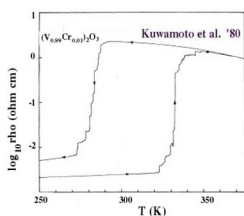
Hysteresis occurs only when both spinodal lines are crossed.





## More with hysteresis: Avalanches

### Avalanches, Barkhausen noise





## *More with hysteresis: Avalanches*

- Requires a theory that include domains and pinning disorder.
- Random field Ising model: RFIM (J. P. Sethna, K. A. Dahmen, E. W. Carlson).

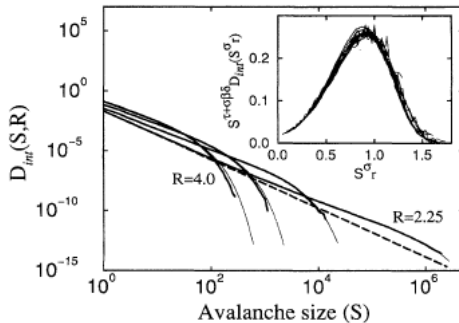
$$H = -J \sum_{ij} S_i S_j - \sum_i (h + h_i) S_i \quad (1)$$

- RFIM can model for Barkhausen noise, for avalanches in MIT system we need a mapping or connection.



## *Barkhausen noise: RFIM*

Olga, Dahmen, Sethna, PRL 75, 4528 ('95)





## *Barkhausen noise: RFIM*

Olga, Dahmen, Sethna, PRL 75, 4528 ('95)

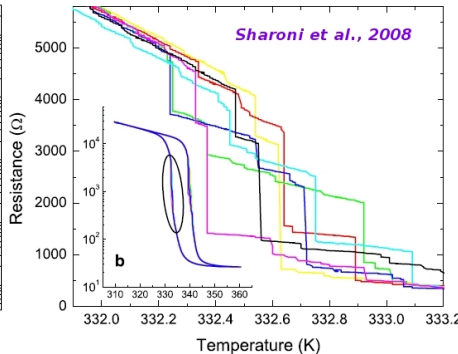
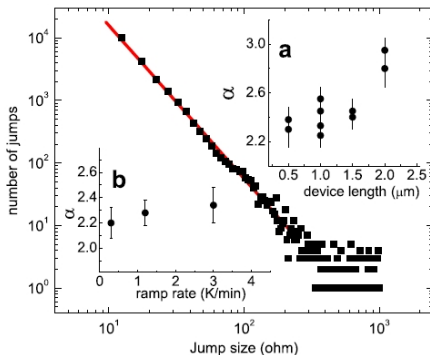
Exponents	Simulation in 3 dimensions	Experiments in 3 dimensions
$\tau$	$1.6 \pm 0.06$	1.74, 1.78, 1.88 [2]; 1.5 $\pm$ 0.5 [4]; 1.33 [10]; 1.5–1.7 [5]
$\tau + \sigma\beta\delta$	$2.03 \pm 0.03$	1.73–2.1 [8]
$(\tau - 1)/\sigma\nu z + 1$	$2.05 \pm 0.12$	1.64, 2.1, 1.82 [2]; 1.7–2 [5]
$(\tau + \sigma\beta\delta - 1)/\sigma\nu z + 1$	$2.81 \pm 0.11$	2.28 [8]
$(3 - \tau)/\sigma\nu z$	$2.46 \pm 0.17$	Around 2 [2,9]
$[3 - (\tau + \sigma\beta\delta)]/\sigma\nu z$	$1.70 \pm 0.10$	1.6 [6,7]; 1.8 [11]
$(\tau - 1)/(2 - \sigma\nu z) + 1$	$1.42 \pm 0.04$	1.44, 1.58, 1.60 [2]



## $VO_2$ experiment

### Jump (avalanche) size distribution

*Sharoni et al., PRL 101, 026404 ('08)*



Power law: exponent,  $\alpha = 2 : 48 \pm 0 : 05$

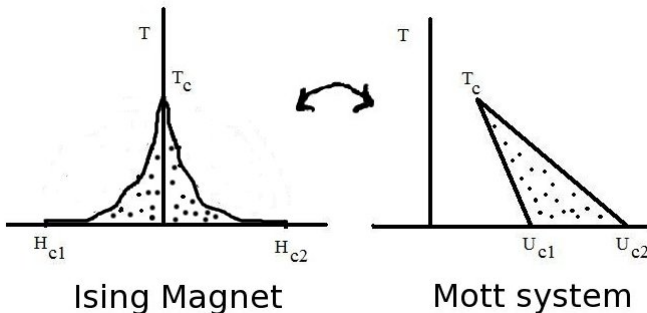


## *Avalanches in MIT: How to model?*

Two ways !

Method 1

A canonical transformation from  $U - T$  to  $h - T$  phase diagram



For simplicity we can think of

$$h(U, T) = a(T)U + bT + c$$



## *Avalanches in MIT: How to model?*

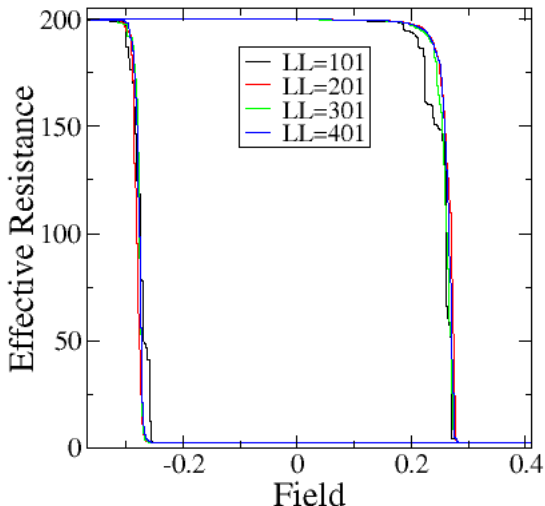
### Method 2: Percolation picture

- A lattice-gas argument for the M-I domains
  - Erica's method: Nematic Stripes  $\rightarrow$  Random field Ising model (RFIM)  $\rightarrow$  Resistor network (RN); *Carlson et al., PRL 96, 097003 ('06)*
  - Similarly Lattice-gas model (M-I)  $\rightarrow$  RFIM  $\rightarrow$  RN
- However, percolation picture does not hold near the critical pt.  
M-I contrast ratio  $\rightarrow 1$
- Need to have tunable resistances in the resistor network.



## *Avalanches in MIT: Finite Size effect*

Avalanche size decreases as system size increases







## Quick Summary

- Hysteresis is common feature in many condensed matter systems.
- Thermal hysteresis in Metal-Insulator systems, e.g.  $V_2O_3$ ,  $VO_2$  can be explained by the phase-diagram geometry.
- RFIM can be suitable for realizing the statistics in the avalanches phenomenon.
- A resistor network model should be connected to the RFIM for the MIT systems.
- The contrast ratio between metallic and insulating resistances should be tuned as the critical point is approached.



## *Acknowledgement*

I wish to thank the Department of Science and Technology (DST), Govt. of India, for supporting my research work and to The Institute for Complex Adaptive Matter (ICAM-I2CAM), 1 Shields Avenue, Davis, CA 95616, and The National Science Foundation (NSF), 4201 Wilson Boulevard, Arlington, VA 22230 (grant number DMR-0844115) for supporting my travel.

Thanks for your kind attention !

