

WHY LHC?

D. P. ROY

Homi Bhabha Centre for Science Education
Tata Institute of Fundamental Research
Mumbai, India

Contents

- Basic Constituents of Matter and their Interactions : Matter Fermions and Gauge Bosons (Std Model)
 - High Energy Colliders
 - Discovery of Std Model Particles at Colliders
 - Higgs Mechanism : Higgs Search at LHC
 - Supersymmetry : SUSY Search at LHC
- (Natural units: $\hbar & c = 1 \Rightarrow m = mc^2$, $m_p \approx 1 \text{ GeV}$)

Basic Constituents of Matter

Mass (GeV)

Fermions (Spin = 1/2 \hbar)

Leptons v_e v_μ v_τ 0

e .0005 μ 0.1 τ 1.8 -1

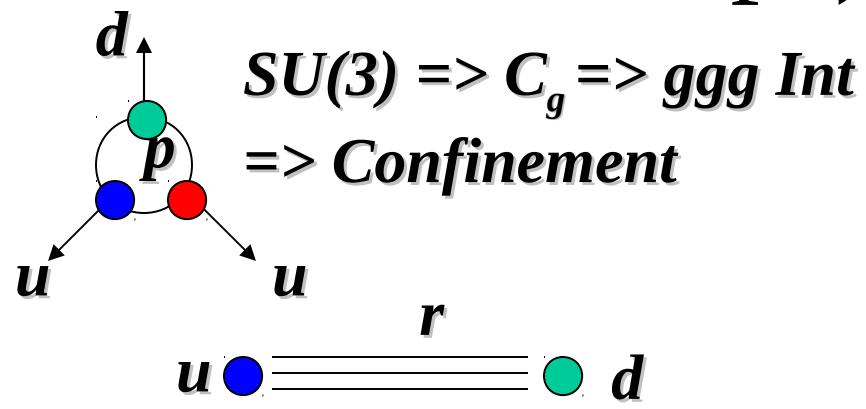
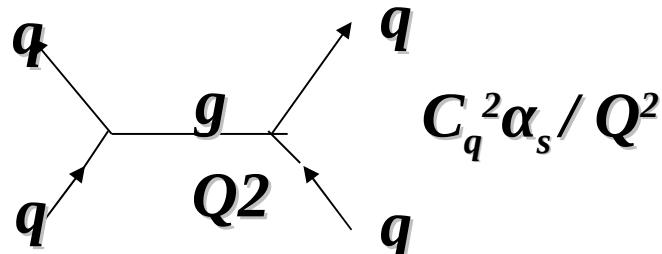
Quarks u 0.3 c 1.5 t 175 2/3 d 0.3 s 0.5
b 5 -1/3

For each Pair : $\Delta e = 1 \Rightarrow$ Weak Int

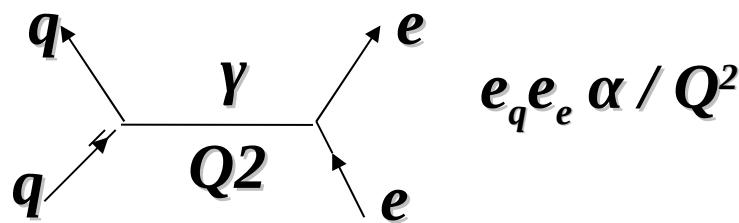
Quarks also carry Colour Charge (C) \Rightarrow Strong Int

Basic Ints (Gauge Bosons & Groups)

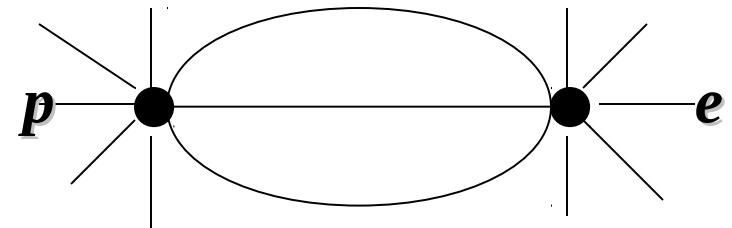
1. Strong Int (QCD) : $SU(3)$



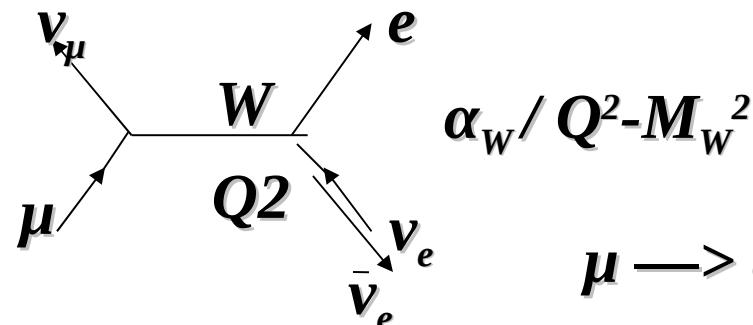
2. E.M. Int (QED) : $U(1)$



$$F = \text{Constant} \Rightarrow V \propto r$$



3. Weak Int : $SU(2)$



$$F = \alpha/r^2 \Rightarrow V = \alpha/r$$

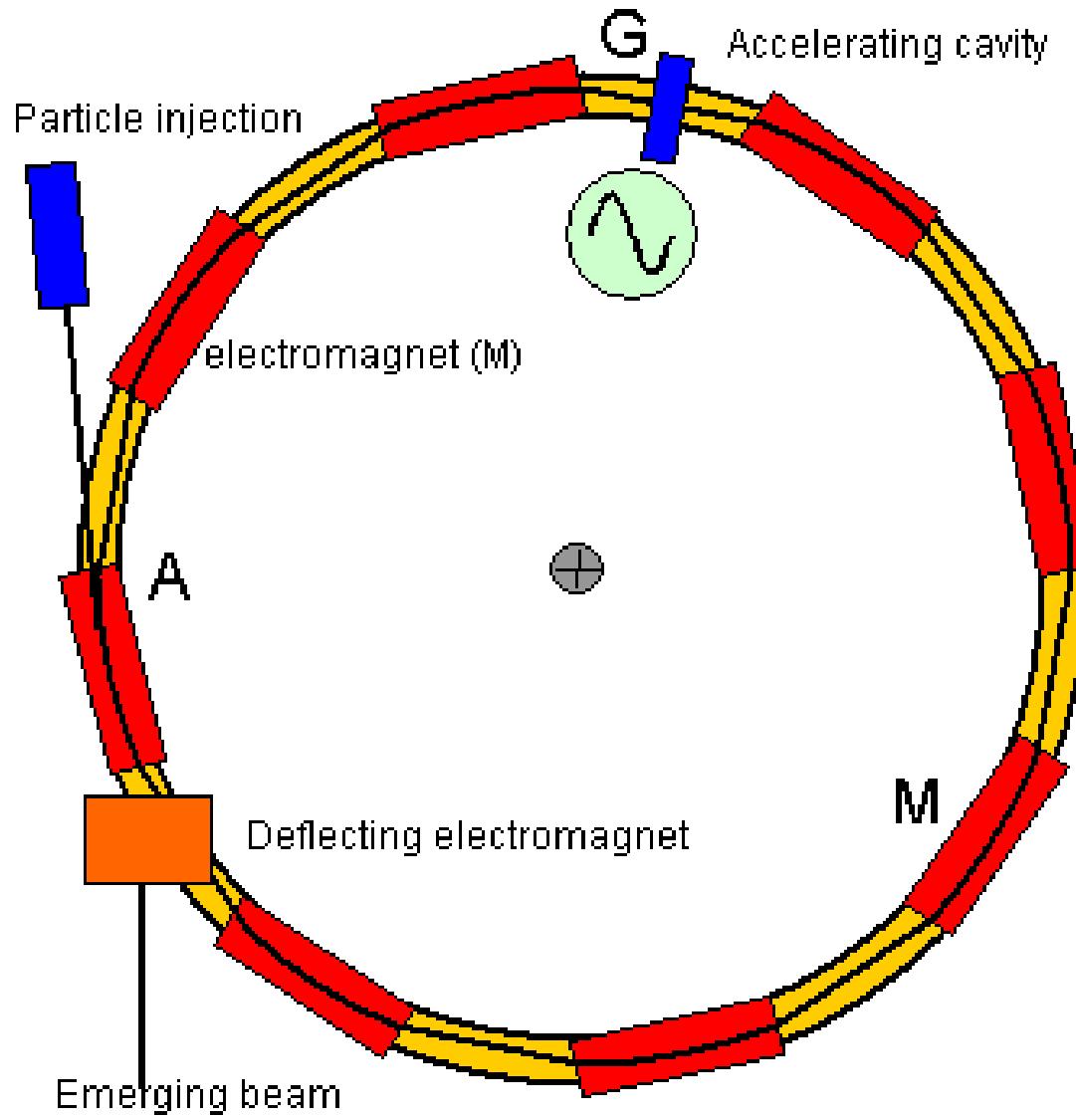
$\Rightarrow V = (\alpha/r) \cdot \exp(-r/M_w)$

$\mu \rightarrow e v_\mu \bar{v}_e : Q^2 \ll M_w^2 \Rightarrow DA_\mu = \alpha_w / M_w^2$

$SU(2) \times U(1)$ EW Th (GSW) $\Rightarrow \alpha_w = \alpha / \sin^2 \theta_w \approx 4\alpha$

$\Rightarrow DA\mu \approx 4\alpha / M_w^2 \Rightarrow M_w = 80 \text{ GeV}, M_z = 91 \text{ GeV}$

- p (uud), n (udd), e \Rightarrow All the Visible Matter
- Heavier Leptons & Quarks Decay by Weak Int
- They can be Observed in Accelerator or Cosmic ray
- Cosmic ray Observation of μ and k ($s\bar{u}$) in 1947
- ν s are stable but very hard to Observe \leq Weak Int
- ν_e Observed in Atomic Reactor Expt in 1956
- ν_μ in BNL PS in 1962 (KGF Cosmic ray in 1965)
- e^+e^- Collider : c (1974), τ (1975), b (1977), g (1979)
- pp Collider : W & Z (1983), t (1995), ν_τ (2000) FT

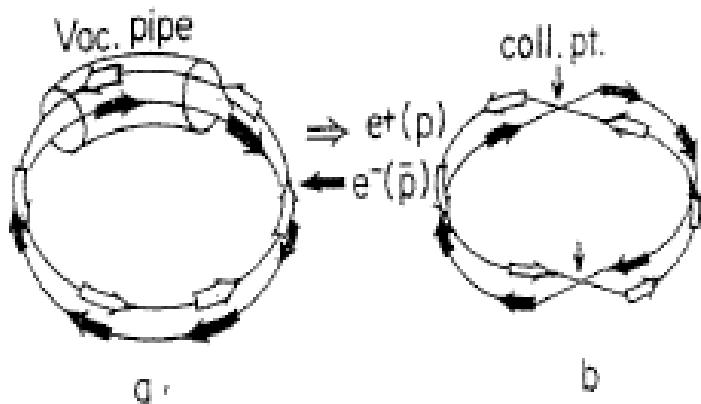


$$eBv = mv^2/R$$
$$eB = mv/R$$

Figure 1

e^+e^- (p^-p) Collider

Acceleration
Mode



Collision
Mode

Advantage of Collider over Fixed Target Accelerator

$$\xrightarrow{E} \xleftarrow{E}$$

$$\sqrt{s} = 2E$$

$$\xrightarrow{E'} \text{[Target Block]}$$

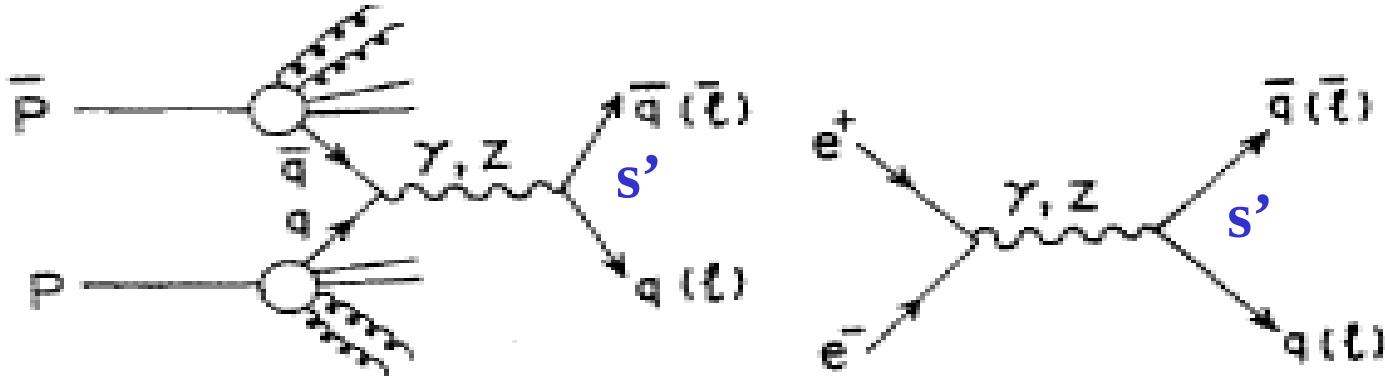
$$\sqrt{s} = \sqrt{2mE'}$$

Tevatron $p\bar{p}$ Collider

$$\sqrt{s} = 2 \times \frac{1000 \text{ GeV}}{1 \text{ TeV}}$$

$$E_{Equiv} = \frac{s}{2m_p} = \frac{(2000 \text{ GeV})^2}{2 \times 1 \text{ GeV}} = 2000,000 \text{ GeV}$$

pp Collider vs e^+e^- Collider



$$s' \sim \langle x_q \rangle s_{p\bar{p}}$$

$$1/6$$

$$s' = s_{e^+e^-}$$

$$\text{Same } \sqrt{s'} \Rightarrow \sqrt{s_{p\bar{p}}} \approx 6\sqrt{s_{e^+e^-}} : \Delta E_{Sync} \approx \frac{4\pi \cdot e^2 E^4}{3m_e^4 \rho}$$

$$\sqrt{s'} \approx M_Z \approx 100 \text{ GeV} : \sqrt{s_{p\bar{p}}} \approx 600 \text{ GeV}, \sqrt{s_{e^+e^-}} \approx 100 \text{ GeV}$$

CERN: $p\bar{p}$ Coll. ($\rho = 1 \text{ km}$), LEP-I ($\rho = 5 \text{ km}$)

COST: (200 + 100) million\$, 1 billion\$

Precision: Tune e^+e^- Energy = $M_Z \Rightarrow$ Higher Rate & Better Mass Res

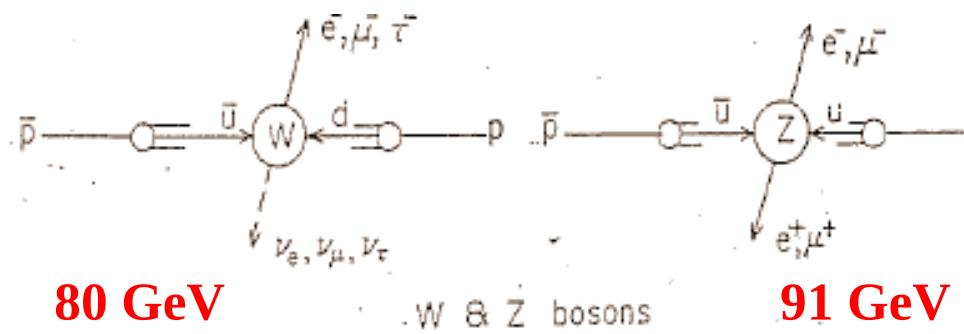
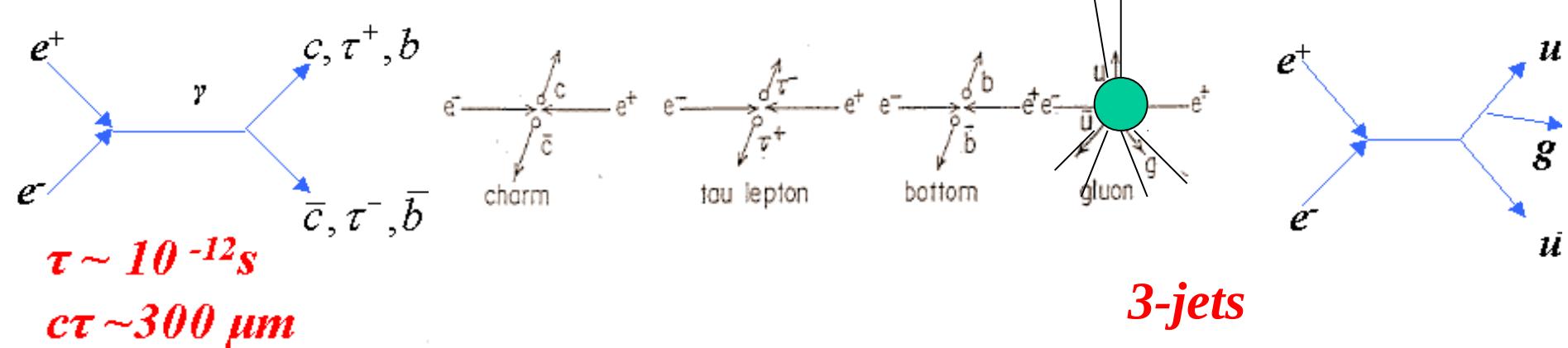
Z events/yr : LEP-I $\sim 10^6$, CERN $p\bar{p}$ Coll. $\sim 10^{1-2}$

Signal : Clean , Dirty (Debris from Spectator q & g)

Past, Present & Proposed Colliders

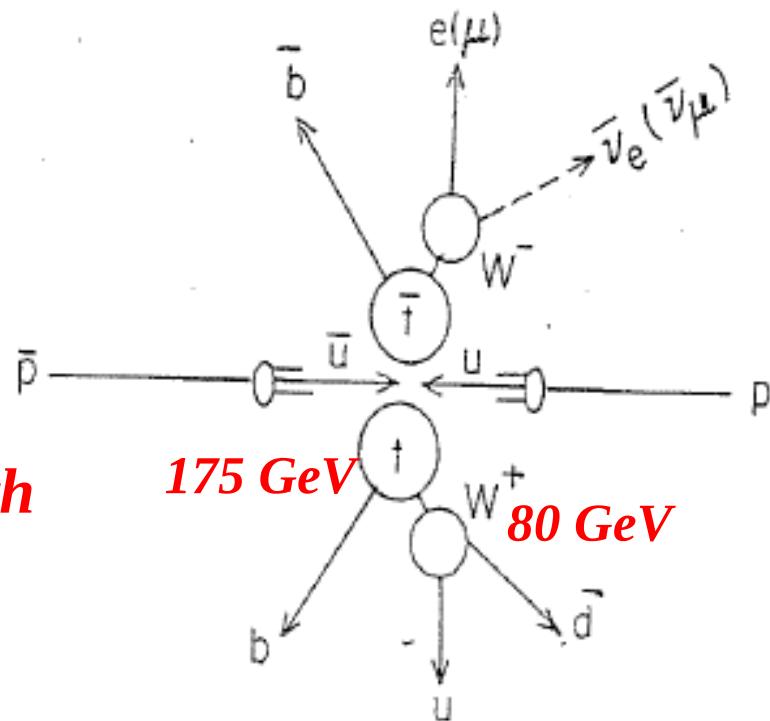
Period	Machine	Location	Beam	Energy(GeV)	Radius	Highlight
70's	SPEAR	Stanford	e ⁺ e ⁻	3 + 3		charm , τ
	DORIS	Hamburg		5 + 5		bottom
	CESR	Cornell		8 + 8	125 m	bottom
	PEP	Stanford		18+18		
	PETRA	Hamburg		22+22	300 m	gluon
80's	TRISTAN	Japan	e ⁺ e ⁻	30+30		
	SPPS	CERN	p p	300+300	1 km	W,Z boson
90's	Tevatron	Fermilab	p p	1000+1000		Top
	SLC	Stanford	e ⁺ e ⁻	50+50		Z
	LEP-I	CERN		50+50	5 km	Z
	(LEP-II)			100+100		W
	HERA	Hamburg		30+800		
			e p			
2009	LHC	CERN	p p	7000+7000	5 km	Higgs,SUSY
2???	ILC	???	e ⁺ e ⁻	500+500		

$$\Leftrightarrow \mathbf{k}_T \sim \hbar/1\text{fm} \sim 0.2 \text{ GeV}$$



**Hard e/μ ($p_T \sim 40 \text{ GeV}$)
with apparent p_T imbalance
(missing - p_T)**

**Hard back-to-back
 $e^+e^-/\mu^+\mu^-$
($p_T \sim 45 \text{ GeV}$)**



*Hard Isolated e / μ with
3 - 4 Hard jets*

175 GeV 80 GeV

Top pair production

Godbole, Pakvasa & Roy, Phys. Rev. Lett. 50, 1539 (1983)

Gupta & Roy, Z. Phys. C39, 417 (1988)

Mass Problem : Higgs Mechanism

How to give mass to the $SU(2)$ Gauge Bosons w/o breaking Gauge Sym of the L ? For simplicity look at the $U(1)$ Gauge Th (EM Int).

$$L_{EM} = (\partial_\mu - ieA_\mu) \overline{\phi^*} (\partial_\mu + ieA_\mu) \overline{\phi} - [\mu^2 \phi^* \phi + \lambda(\phi^* \phi)^2] - \frac{1}{4} F_{\mu\nu} F_{\mu\nu}$$

$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu \Leftrightarrow E_{1..3}, B_{1..3}$$

~~$-M^2 A_\mu A_\mu$~~

$$GaugeTr : \phi \rightarrow e^{i\alpha(x)} \phi, A_\mu \rightarrow A_\mu + \frac{1}{e} \partial_\mu \alpha(x)$$

$$V = \mu^2 \phi^* \phi + \lambda(\phi^* \phi)^2$$

$$\mu^2 < 0$$

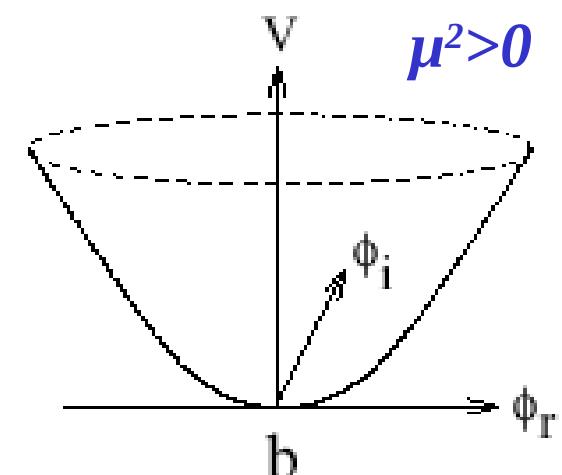
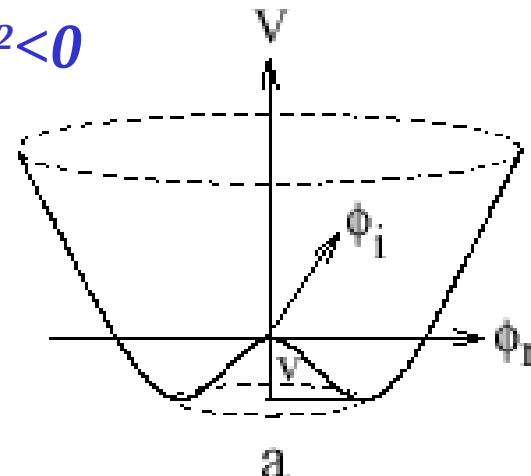
$$v = \sqrt{-\mu^2 / \lambda} \quad \text{vev}$$

$$\Rightarrow \phi_r = v + h(x)$$

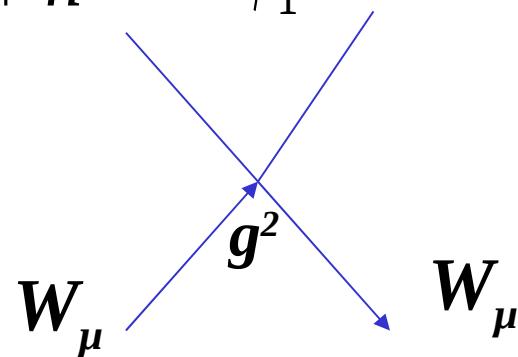
$$\Rightarrow M = ev$$

$$m_q = y_q v$$

$$\Rightarrow m_h = \sqrt{-\mu^2} = \sqrt{\lambda} v = \sqrt{\lambda} / e \times M \sim 10^2 \text{ GeV}$$

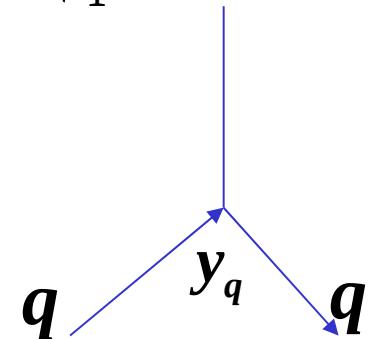


$$\phi_1^0 = v + h$$



$$\phi_1^0 = v + h$$

$$\phi_1^0 = v + h$$



$$g^2 v^2 W_\mu W_\mu + g^2 v h W_\mu W_\mu$$

M_W^2 $g M_W$

$$y_q v \bar{q}q + y_q h \bar{q}q$$

m_q m_q/v

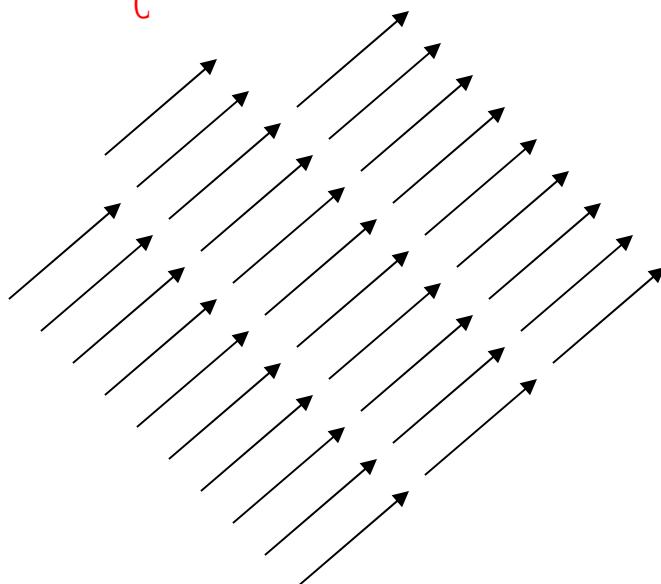
Higgs couplings to Particles is Proportional to their Mass
=>Most Important Channels for Higgs Search are the Heavy Pairs: h (WW , ZZ , tt , bb , $\tau\tau$) & H^\pm (tb , $\tau\nu$)

τ Polarization Effect <= Roy, Phys.Lett.B459, 607 (1999)

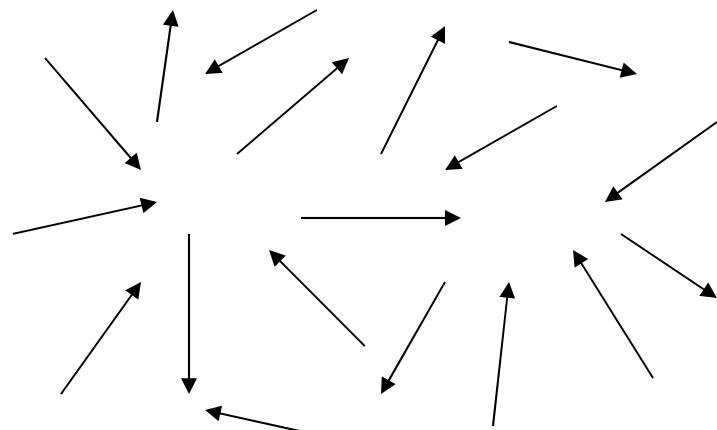
Ferromagnetism

(Spontaneous Symmetry Breaking)

$T < T_c$

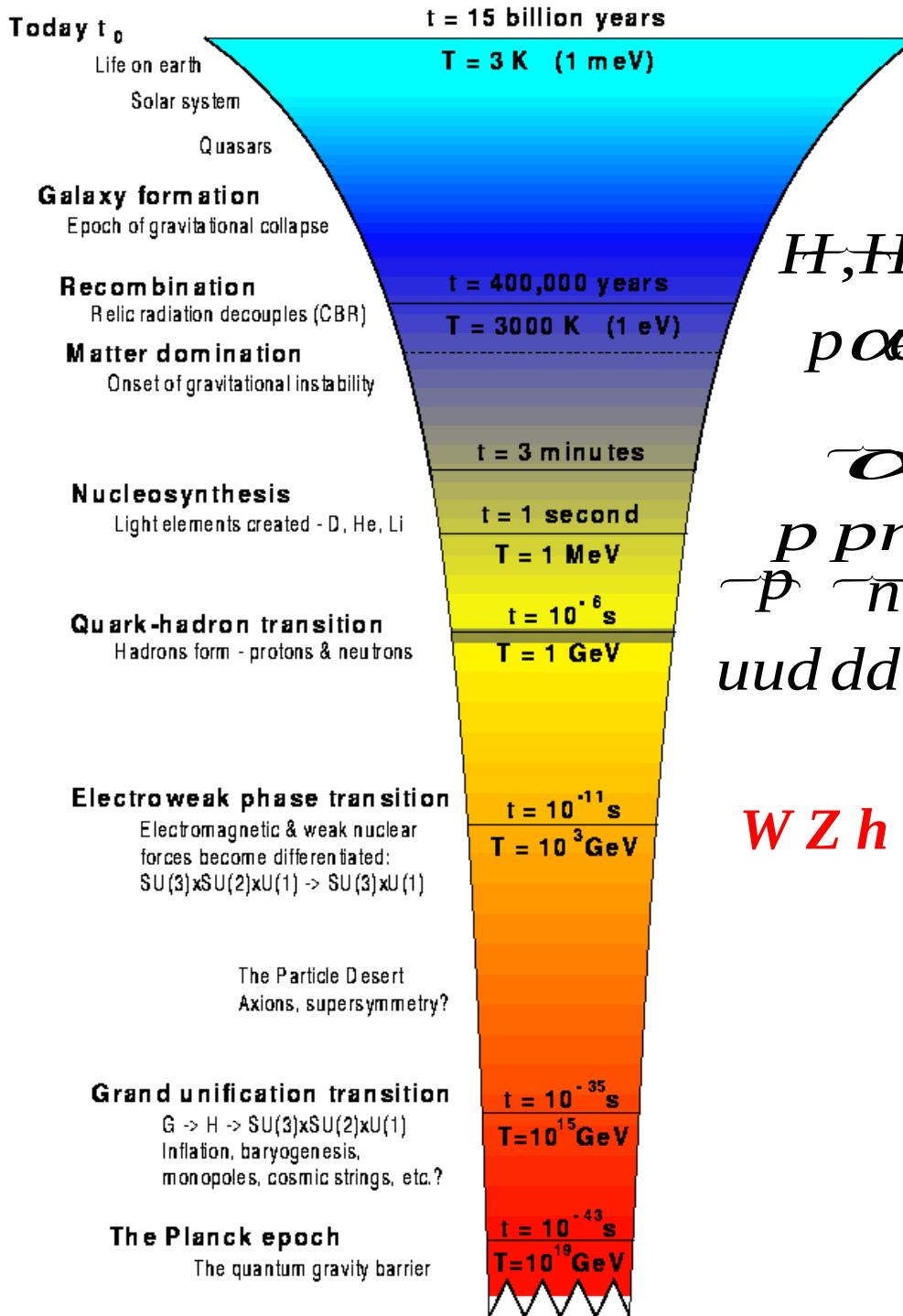


$T > T_c$



Rotational Symmetry Broken

Rotational Symmetry



H, He

$p \alpha e \nu \gamma$

α

$p p n e \nu \gamma$
 $\bar{p} \bar{n}$

$u u d d \bar{u} \bar{d} e \nu \gamma$

$W Z h t b e \nu \gamma$

Hierarchy Problem: Supersymmetry (SUSY) Solution

How to control Higgs Mass $m_h \sim M_W \sim 10^2 \text{ GeV}$?

$$\frac{1}{(k^2 - m^2)} \quad \text{Diagram: } \begin{array}{c} \text{circle} \\ k \end{array} \rightarrow \lambda \int \frac{k^3 dk}{(k^2 - m^2)} \propto k^2 \triangleright \infty \quad - h \quad \tilde{h}$$

Without a protecting symmetry scalar mass gets quad. div. quantum corr. $\Rightarrow m_h \rightarrow \infty(M_{\text{GUT}}; M_{\text{Planck}})$

SUSY: $10^{16} \quad 10^{19} \text{ GeV}$

fermions<=>bosons

	s	s	s
q, l	$1/2$	γ, g, W, Z	1
\tilde{q}, \tilde{l}	0	$\tilde{\gamma}, \tilde{g}, \tilde{W}, \tilde{Z}$	$1/2$

=> Superparticles mass $\sim 10^2 \text{ GeV}$

$$R = (-1)^{3B + L - 2s}$$

	$+1$
$1/2$	-1

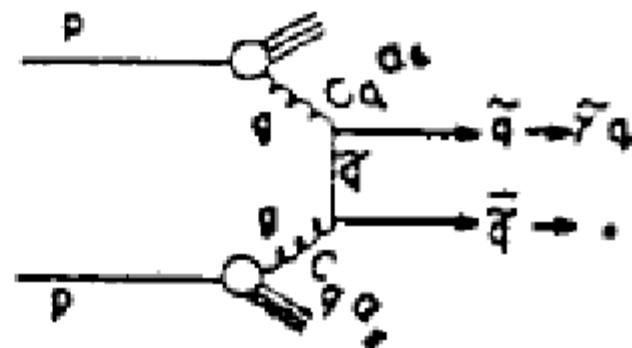
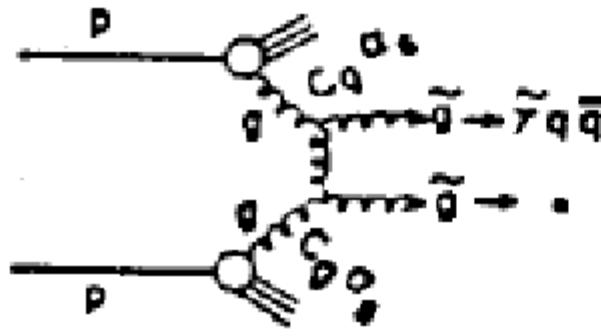
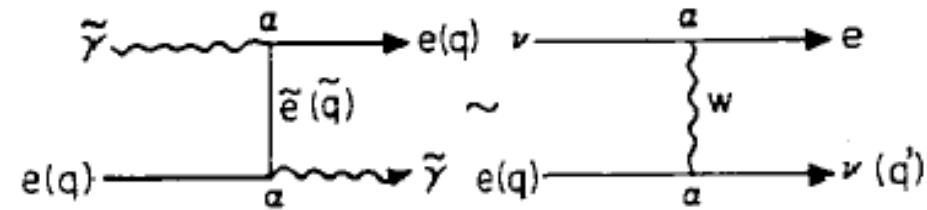
R Cons. => Pair-prod of SP & Stable LSP (Cold dark matter)

LSP $\tilde{\chi}$ Weakly Int Massive Particle (WIMP)

$$M_{\tilde{e}, \tilde{q}} \approx M_W$$

=> LSP escapes detection like ν

=> Apparent imbalance of P_T
(Missing- P_T Signature)



Pair production of Gluinos and Squarks at LHC

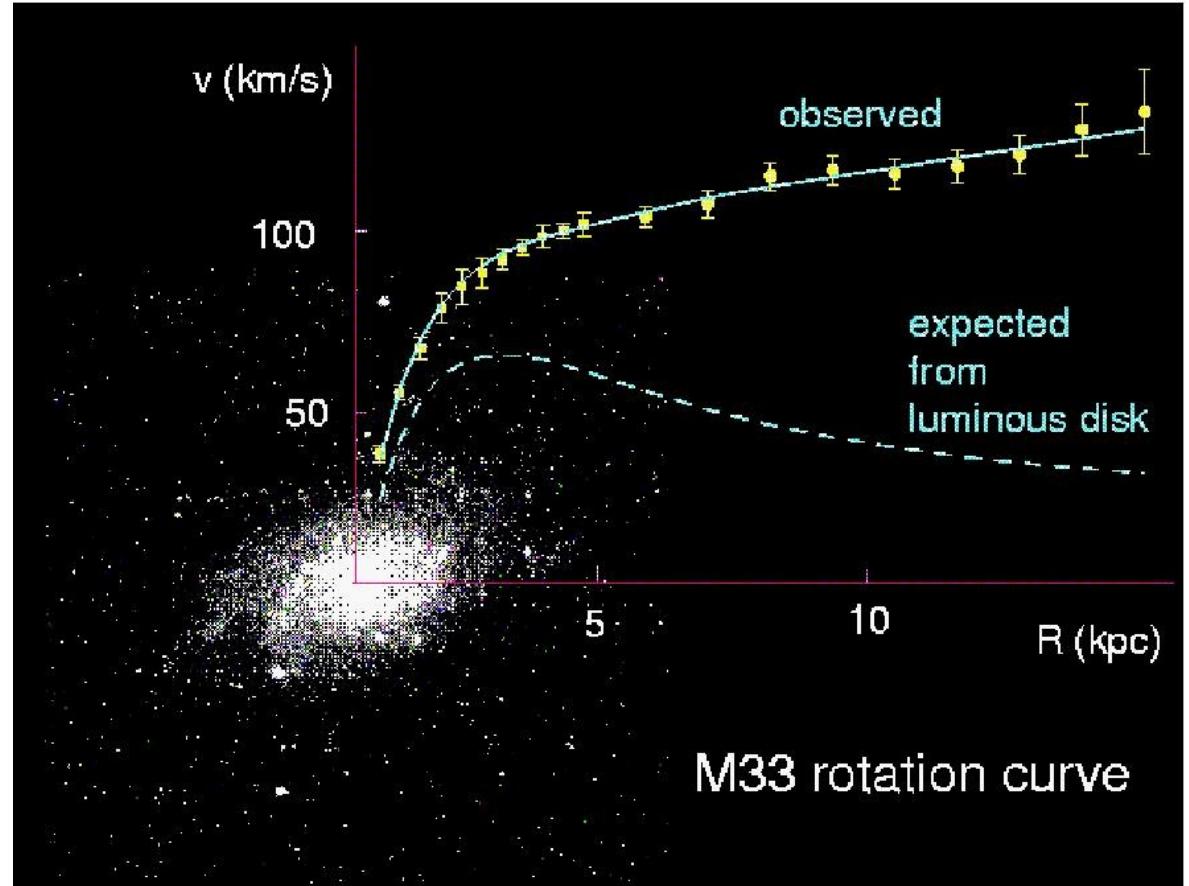
=> Multi-jet plus Missing- P_T Signature for SUSY

Reya & Roy, Phys. Lett. 141B, 442 (1984); Phys. Rev. Lett. 53, 881 (1984)

Conclusion

- Higgs & Superparticles are the minimal set of missing pieces reqd. complete the picture of particle physics (MSSM).
- LHC offers comprehensive Higgs and Superparticle search up to $M_{H,SUSY} = 1000\text{GeV}$.
- It will either complete the picture a la MSSM or provide valuable clue for an alternative picture :
Little Higgs, Extra Dim, ETC...↓
- LSP is leading candidate for the cosmic dark matter ~ 10 times the baryonic matter of the Universe.

Rotation curve of nearby dwarf spiral galaxy M33, superimposed on its optical image



M33 rotation curve

$$\frac{v^2}{R} = G \frac{M(R)}{R^2} \implies v \propto \sqrt{\frac{M(R)}{R}}$$

