

# Searches for BSM

Shrihari Gopalakrishna



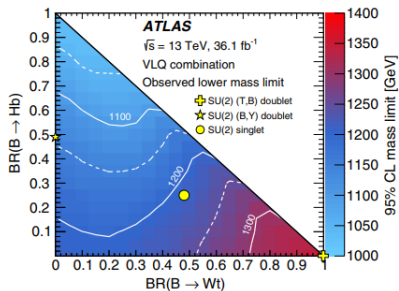
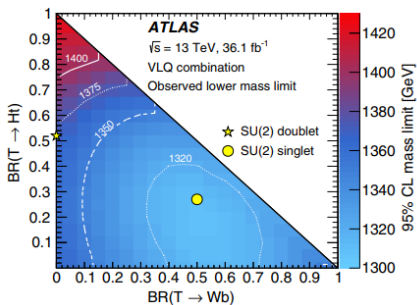
Institute of Mathematical Sciences (IMSc), Chennai

*WHEPP-2019*

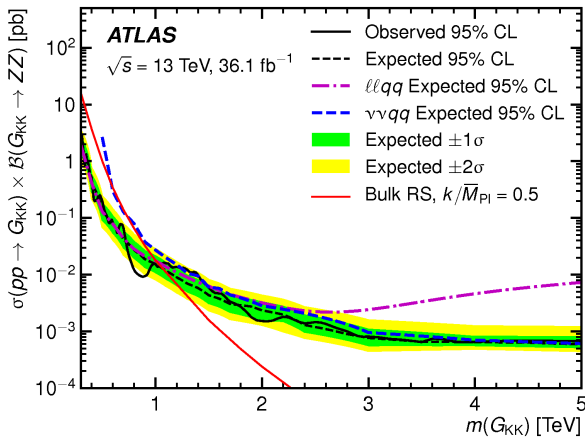
*IIT Guwahati, Dec 2019*

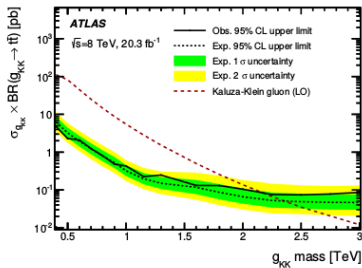
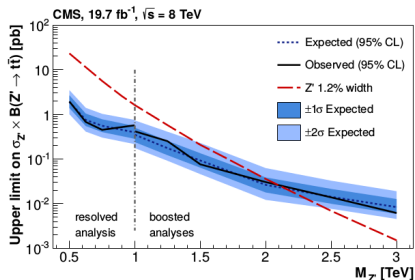
# Talk Outline

- General aspects of searches for BSM

Vector-like fermion ( $t', b'$ ) search

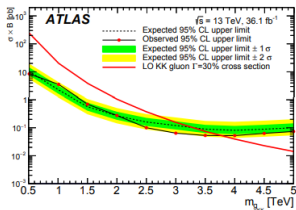
[ATLAS: 1808.02343; PRL 2018]

graviton<sup>(1)</sup>  $\rightarrow ZZ$  search

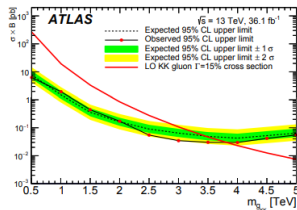
gluon<sup>(1)</sup>  $\rightarrow t\bar{t}$  search (8TeV)(b)  $g_{KK}$ , resolved and boosted combination.Limit:  $M_{KK} > 2.2$  TeV @ 95% CL

[ATLAS 1505.07018; CMS 1309.2030]

# gluon<sup>(1)</sup> $\rightarrow t\bar{t}$ search (13TeV)



(a) 30% width

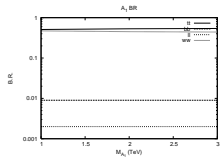
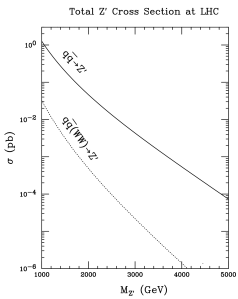
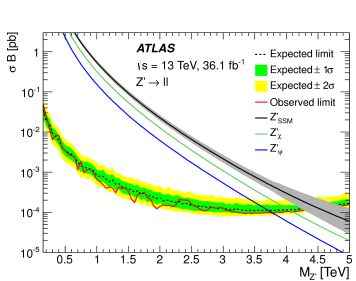


(b) 15% width

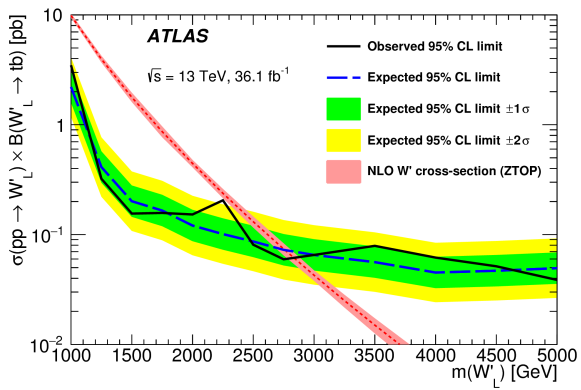
Figure 17: The observed and expected cross-section 95% CL upper limits on the  $g_{KK}$  signal for resonance widths of (a) 30% and (b) 15%. The theoretical predictions for the production cross-section times branching ratio of  $g_{KK} \rightarrow t\bar{t}$  at the corresponding masses are also shown.

[ATLAS 1804.10823 [hep-ex]]

$$Z' \rightarrow l^+ l^-$$



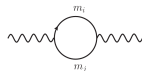
[Agashe, Davoudiasl, SG, Han, Huang, Perez, Si, Soni - 0709.0007 [hep-ph]]





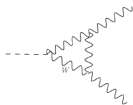
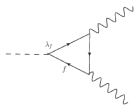
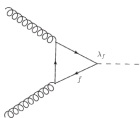
# EWPrecision + Higgs Observables

Precision electroweak observables ( $S, T, U$ )



Modifications to  $hgg$ ,  $h\gamma\gamma$  couplings:

$\sigma(gg \rightarrow h)$        $\Gamma(h \rightarrow \gamma\gamma)$

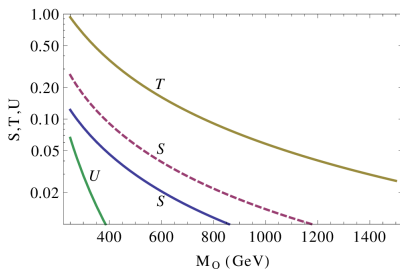
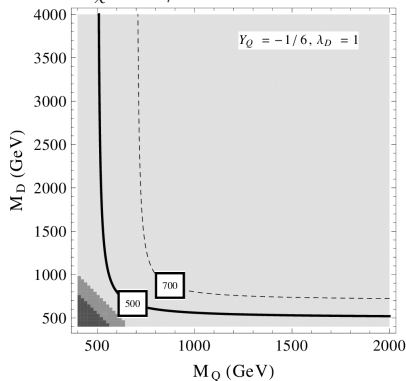


We compute ratios  $\frac{\Gamma_{h \rightarrow gg}}{SM}$ ,  $\frac{\Gamma_{h \rightarrow \gamma\gamma}}{SM}$

using leading-order expressions

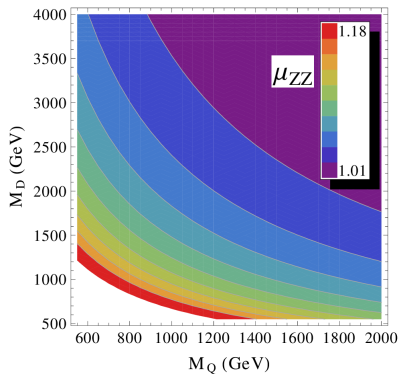
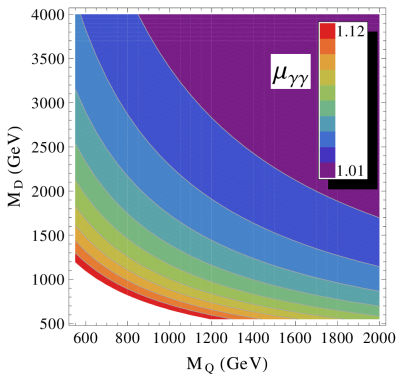
QCD corrections to ratios small: [Furlan '11] [Gori, Low '13]

$$\mu_{\gamma\gamma}^{VBF} \approx \frac{\Gamma_{\gamma\gamma}}{\Gamma_{\gamma\gamma}^{SM}}; \quad \mu_{ZZ}^{ggh} \approx \frac{\Gamma_{gg}}{\Gamma_{gg}^{SM}}; \quad \mu_{\gamma\gamma}^{ggh} \approx \frac{\Gamma_{gg}}{\Gamma_{gg}^{SM}} \frac{\Gamma_{\gamma\gamma}}{\Gamma_{\gamma\gamma}^{SM}}; \quad \frac{\mu_{\gamma\gamma}^{ggh}}{\mu_{ZZ}^{ggh}} \approx \frac{\Gamma_{\gamma\gamma}}{\Gamma_{\gamma\gamma}^{SM}} \approx \mu_{\gamma\gamma}^{VBF}$$

$2\bar{2} + 1\bar{1}$  model $Q + U$  model (ST Model like) : MVQD Model with  $Y_\chi = -1/6$  $\lambda_D = 1, M_D = M_Q, Y_Q = (1/6, -1/6)$  (solid, dashed)

[S.Ellis, R.Godbole, SG, J.Wells; 1404.4398, JHEP 2014]

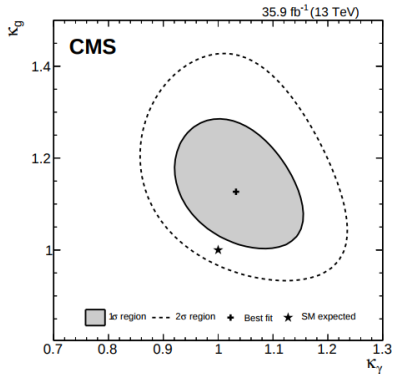
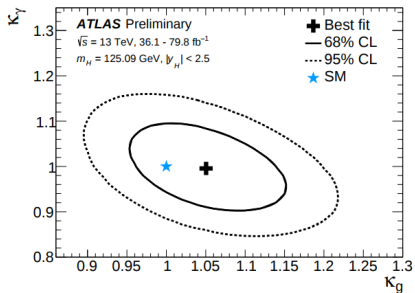
# $Q + U$ model



[Q+U model from MVQD model with  $Y_\chi = -1/6$ ]

[S.Ellis, R.Godbole, SG, J.Wells; 1404.4398, JHEP 2014]

# LHC constraints on Higgs couplings



[ATLAS-CONF-2018-31] [CMS-HIG-17-031]

# BACKUP SLIDES

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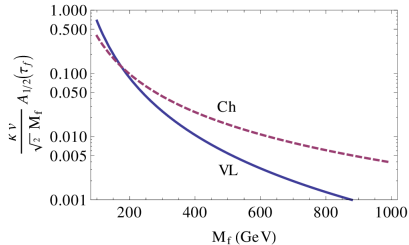
# BSM Vector-like Fermions (VLF)

Vector-like fermions have both  $L$  and  $R$  chiralities charged under a gauge-group.  
This allows a bare mass term.

- VLFs appear in many BSM extensions  
(Eg: composite-Higgs theories, Extra-dimensional theories)
  - ▶ they are sometimes the lightest BSM states
- We study VLF effects on Higgs vacuum stability
  - ▶ constraint on parameter space
    - but any other new states will alter conclusions!

# Vector-like fermion (VLF) decoupling

- VLF has independent source of mass  $M$  (not given by  $m = \lambda v$ )
  - ▶ Can make  $M$  arbitrarily large
    - Yukawa coupling can be small; so perturbative
  - ▶ Nice decoupling behavior :  $S, T, U, h \rightarrow \gamma\gamma, gg \rightarrow h, \dots$ 
    - For instance  $h\gamma\gamma, ggh$  couplings



# VLF signatures

## Observables

- Precision Electroweak Probes
- LHC signals
  - ▶ Direct:  $b' \rightarrow tW, bZ$ ;  $t' \rightarrow bW, tZ, th$ ;  $\chi \rightarrow tW$
  - ▶ Indirect: Higgs coupling modifications
- FCNC probes
- Vacuum stability implications



# Precision Electroweak Constraints

Precision Electroweak Constraints ( $S, T, Zb\bar{b}$ )  
(perturbatively calculable on the warped side)



- Bulk gauge symm -  $SU(2)_L \times U(1)$  (SM  $\psi$ , H on TeV Brane)
- T parameter  $\sim (\frac{v}{M_{KK}})^2 (k\pi R)$  [Csaki, Erlich, Terning 02]
  - ▶ S parameter also  $(k\pi R)$  enhanced
- AdS bulk gauge symm  $SU(2)_R \Leftrightarrow$  CFT Custodial Symm [Agashe, Delgado, May, Sundrum 03]
  - ▶ T parameter - Protected; S parameter -  $\frac{1}{k\pi R}$  for light bulk fermions
  - ▶ **Implies heavy vector bosons:**  $W'_\mu, Z'_\mu, \dots$
  - ▶ Problem:  $Zb\bar{b}$  shifted
- 3rd gen quarks (2,2) [Agashe, Contino, DaRold, Pomarol 06]
  - ▶  $Zb\bar{b}$  coupling - Protected
  - ▶ Precision EW constraints  $\Rightarrow M_{KK} \gtrsim 1.5 - 2.5$  TeV
  - ▶ **Implies top partners:**  $t', b', \chi, \dots$

# Warped Fermions

- SM fermions :  $(+, +)$  BC  $\rightarrow$  zero-mode
- “Exotic” fermions :  $(-, +)$  BC  $\rightarrow$  No zero-mode
  - ▶ 1<sup>st</sup> KK vectorlike fermion

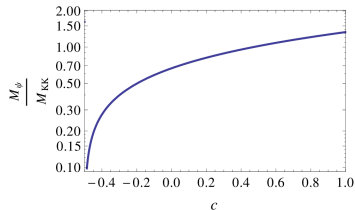
- Typical  $c_{t_R}, c_{t_L}$  :  $(-, +)$  top-partners “light”

$c$  : Fermion bulk mass parameter

[Choi, Kim, 2002] [Agashe, Delgado, May, Sundrum, 03]

[Agashe, Perez, Soni, 04] [Agashe, Servant 04]

- ▶ Look for it at the LHC



[Dennis et al, '07] [Carena et al, '07] [Contino, Servant, '08]

[Atre et al, '09, '11] [Aguilar-Saavedra, '09] [Mrazek, Wulzer, '09]

[SG, Moreau, Singh, '10] [SG, Mandal, Mitra, Tibrewala, '11] [SG, Mandal, Mitra, Moreau : '13]

# Fermion rep : $Zb\bar{b}$ not protected (DT model)

[Agashe, Delgado, May, Sundrum '03]

- Complete  $SU(2)_R$  multiplet

- ▶  $Q_L \equiv (\mathbf{2}, \mathbf{1})_{1/6} = (t_L, b_L)$

- ▶  $\psi_{t_R} \equiv (\mathbf{1}, \mathbf{2})_{1/6} = (t_R, b')$

- ▶  $\psi_{b_R} \equiv (\mathbf{1}, \mathbf{2})_{1/6} = (T, b_R)$

- "Project-out"  $b'$ ,  $T$  zero-modes by  $(-, +)$  B.C.

- New  $\psi_{VL} : b', T$

- $b \leftrightarrow b'$  mixing

- ▶  $Zb\bar{b}$  coupling shifted

- So LEP constraint quite severe

# Fermion rep : $Zb\bar{b}$ protected (ST & TT models)

- $Q_L = (2, 2)_{2/3} = \begin{pmatrix} t_L & \chi \\ b_L & T \end{pmatrix}$

[Agashe, Contino, DaRold, Pomarol '06]

- ▶  $Zb_L\bar{b}_L$  protected by custodial  $SU(2)_{L+R} \otimes P_{LR}$  invariance  
 $W_{t_L b_L}, Z_{t_L t_L}$  not protected, so shifts

Two  $t_R$  possibilities:

① Singlet  $t_R$  (ST Model) :  $(1, 1)_{2/3} = t_R$       New  $\psi_{VL} : \chi, T$

② Triplet  $t_R$  (TT Model) :

$$(1, 3)_{2/3} \oplus (3, 1)_{2/3} = \psi'_{t_R} \oplus \psi''_{t_R} = \begin{pmatrix} \frac{t_R}{\sqrt{2}} & \chi' \\ b' & -\frac{t_R}{\sqrt{2}} \end{pmatrix} \oplus \begin{pmatrix} \frac{t''}{\sqrt{2}} & \chi'' \\ b'' & -\frac{t''}{\sqrt{2}} \end{pmatrix}$$

New  $\psi_{VL} : \chi, T, \chi', b', \chi'', t'', b''$

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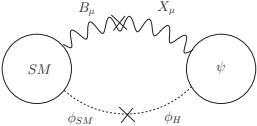
New  $\psi_{VL} : \chi, T, \chi', b', \chi'', t'', b''$

# Hidden sector DM $\psi$

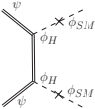
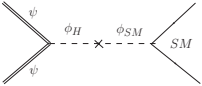
[SG, Lee, Wells 2009]

SM  $\times$   $U(1)_X$  :  $U(1)_X$  sector:  $X_\mu, \Phi_{hid}, \psi$

$$\mathcal{L} \supset -\alpha |H|^2 |\Phi_{hid}|^2 + \frac{\eta}{2} X_{\mu\nu} B^{\mu\nu} - \kappa \phi_{hid} \bar{\psi} \psi$$



Higgs portal DM: Self-annihilation



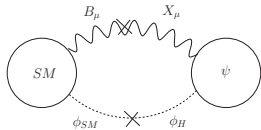
Channels  $\psi\psi \rightarrow b\bar{b}, W^+W^-, ZZ, hh, t\bar{t}$

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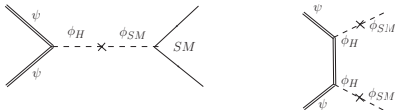
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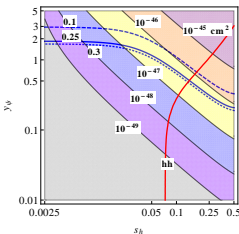
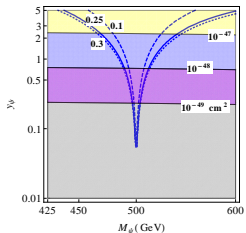
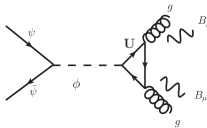
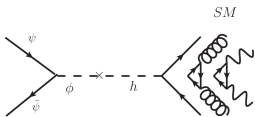


Channels  $\psi\psi \rightarrow b\bar{b}, W^+W^-, ZZ, hh, t\bar{t}$

# Aside: Application to Higgs Portal DM (VLL)

Can also apply to Higgs portal DM case:

[SG. T. Mukherjee: AHEP 2017]



Constraint requires  $s_h \ll 1$ , so vacuum stability constraint is with VLL (DM) effectively coupling with  $\tilde{y} \equiv y_\psi s_h \ll 1$