

Chandrasekhar's contribution to Radiative Transfer: an appreciation

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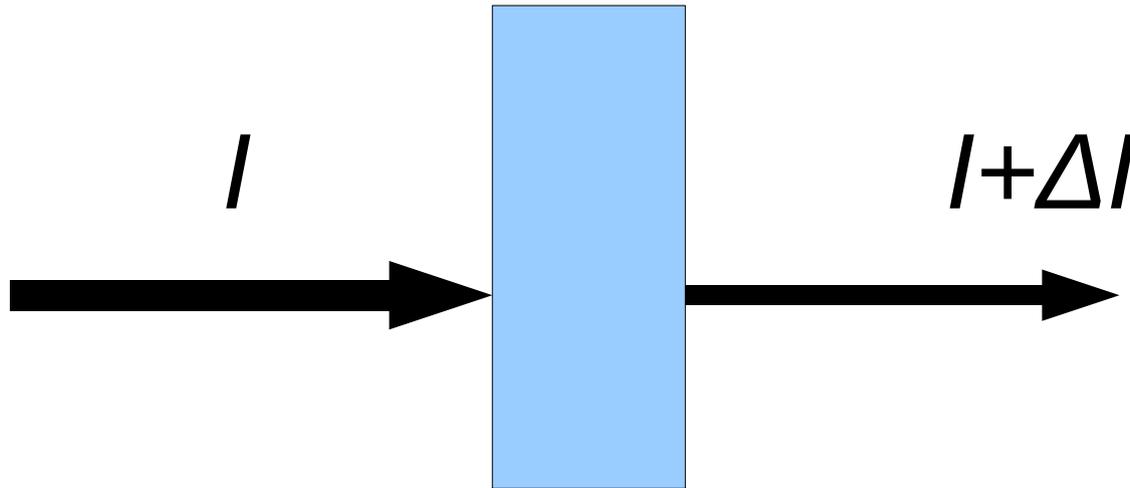
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Plan

- Radiative transfer and how it stood ca. 1944
- A numerical scheme and some analytical fallout
- The principles of invariance
- Blue sky research - polarisation
- The negative ion of hydrogen
- Matters of taste, style, substance

Simple absorption



$$\Delta I / I = -\alpha dl = -d\tau$$

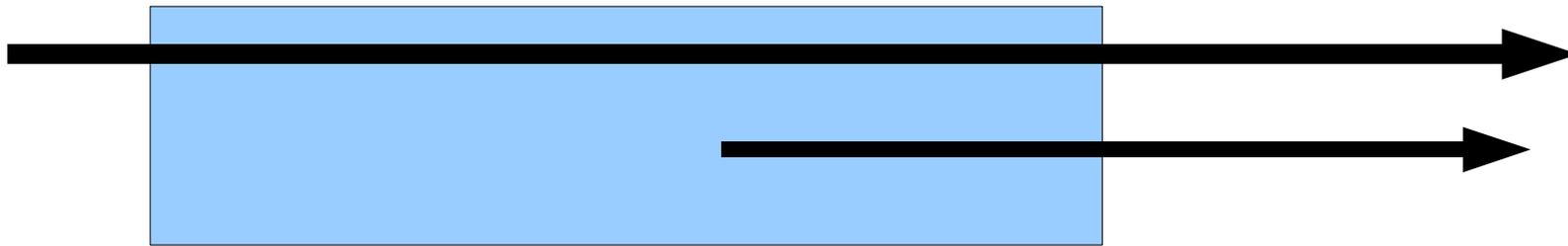
$$I(x) = I(0) \exp\left(-\int_0^x \alpha(y) dy\right) = I(0) \exp(-\tau)$$

Absorption plus emission

$$\Delta I = -\alpha dl I + \epsilon dl; \frac{dI}{dl} = -\alpha I + \epsilon$$

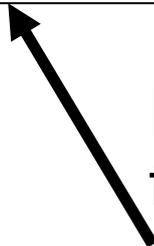
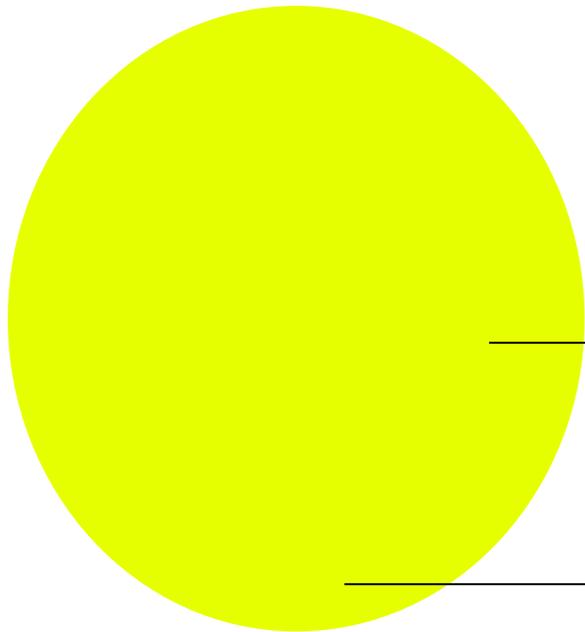
$$I = I(0) \exp(-\tau(0, x))$$

$$+ \int_0^x \epsilon(z) \exp(-\tau(z, x)) dz$$



Note mean free path interpretation – we see down to the 'photosphere' “Optical depth” is the “tau of astrophysics:” distance in units of the mean free path”

Limb darkening



Looking at the edge of the Sun, the surface of optical depth unity is reached at a greater height because of viewing at an angle. This layer is darker

Stellar Atmospheres

- Given temperature as a function of height, the fraction of different ions (Saha) and the detailed behaviour of each for all wavelengths ((Los Alamos tables)
- The reward for modeling and studying stellar spectra is information on chemistry, rotation, magnetic field, turbulence, and even planets!
- In hot stars, radiation is a major player, for self consistency we need “non- LTE “
- Extreme – electron scattering atmosphere

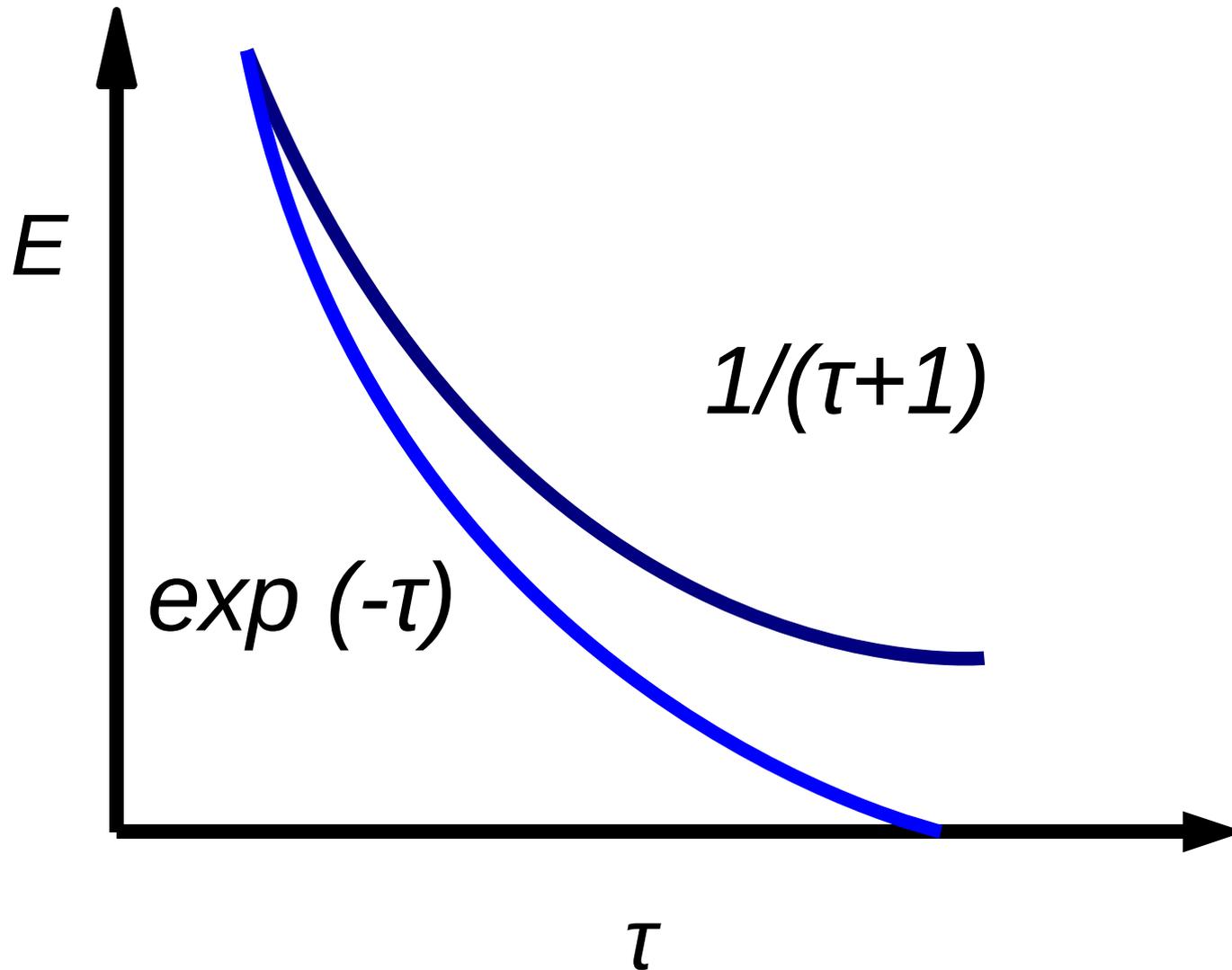
Radiative transfer with scattering

- Earth's atmosphere at visible wavelengths, and those of hot stars (free electrons can't absorb or emit single photons)
- Neutrons in reactors while they are slowing down
- Mathematically more complex and physically more subtle than the pure absorption-emission case

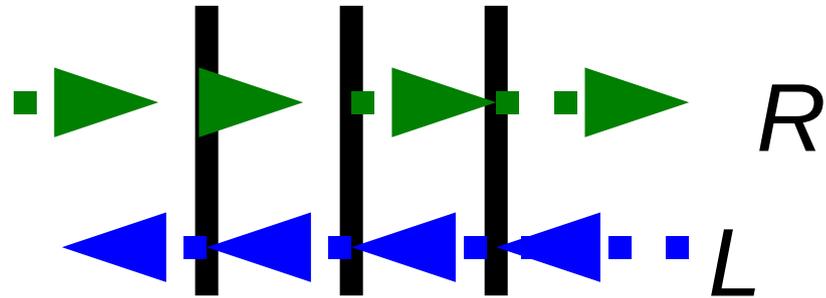
Scattering RT begins at home



Two kinds of attenuation



One-dimensional model with scattering



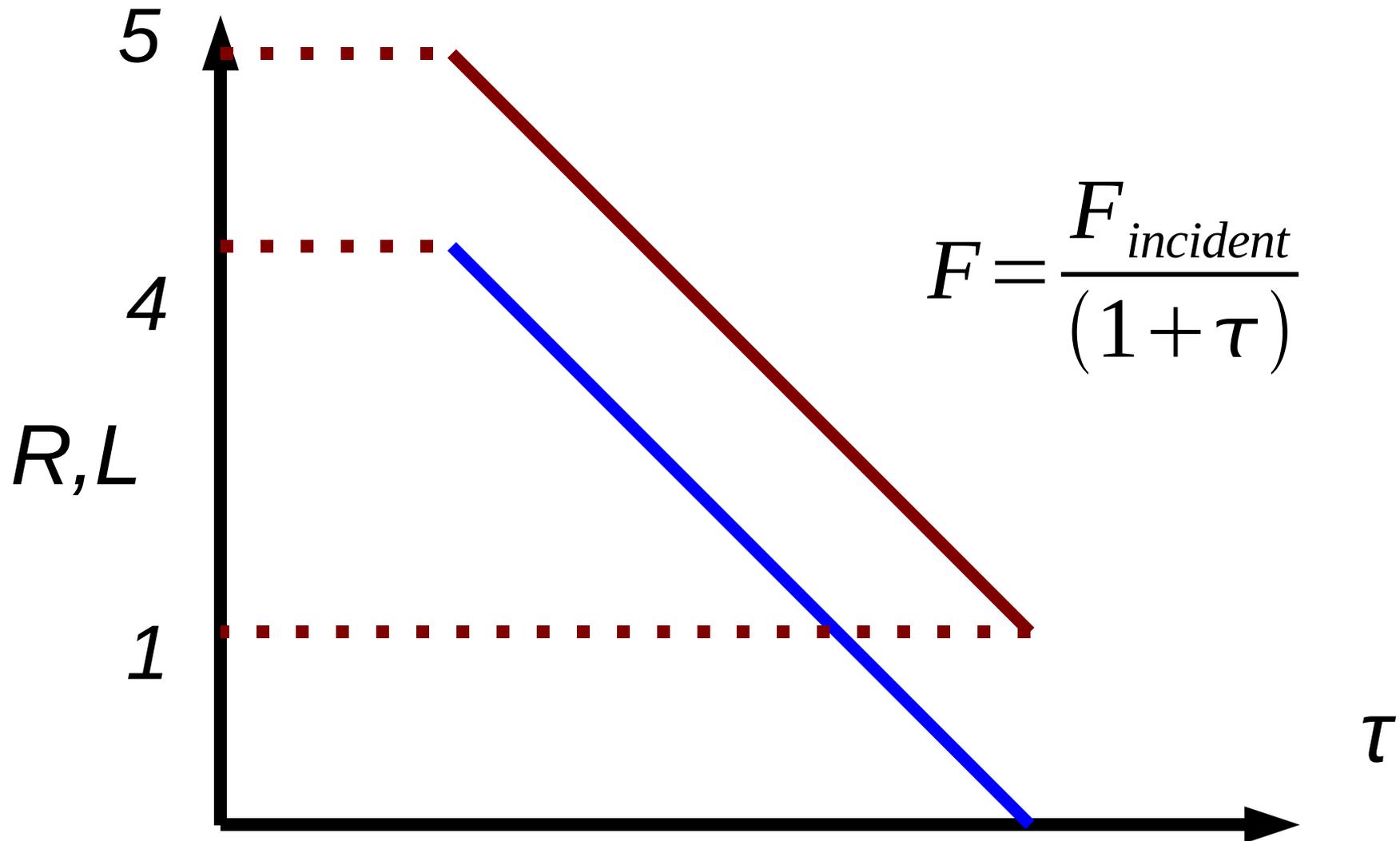
$$dR/dx = -\alpha_s(R) + \alpha_s(L+R)/2$$

$$-dL/dx = -\alpha_s(L) + \alpha_s(L+R)/2$$

$$R - L = F, R + L = E,$$

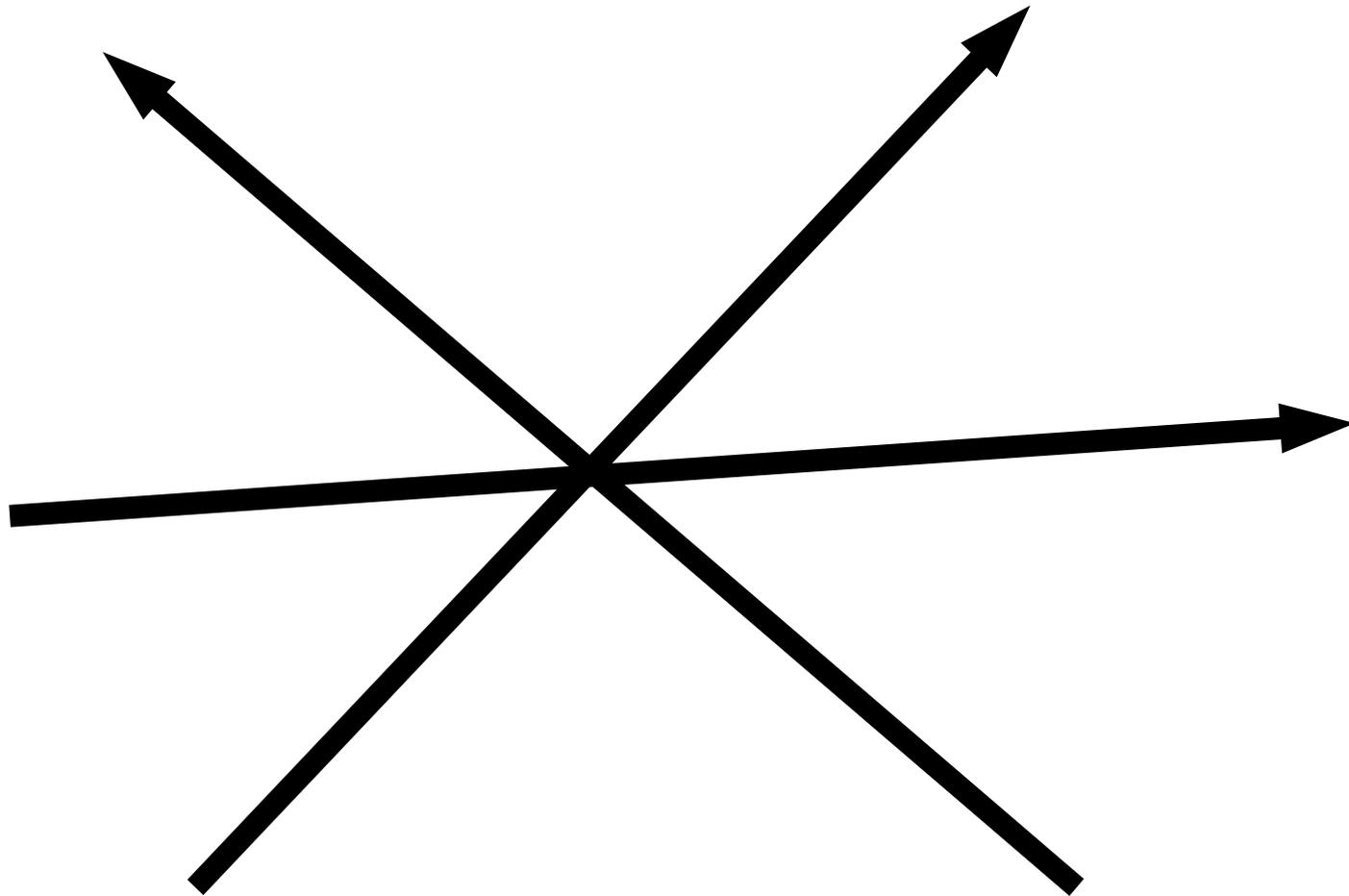
$$dF/dx = 0, dE/d\tau = -F$$

Forward and backward fluxes



Coupling of intensity along different rays by scattering

$$\epsilon(\mu) = \int d\mu' \sigma(\mu, \mu') I(\mu')$$



Coping with scattering

- Equations for moments of the radiation field in the axisymmetric case, e.g

$$\mu \frac{dI}{d\tau} = -I + \frac{1}{2} \int_{-1}^{+1} I d\mu = -I + J$$

$$\int I d\mu = J; \int I \mu d\mu = F; \int I \mu^2 = K, \text{ etc.}$$

- But the equations do not close because of the extra μ on left side
- Eddington cut the Gordian knot by setting $K=J/3$ which is exact for a sphere and hemisphere! Good enough for many people but not for Chandrasekhar

Wiener-Hopf integral equation for the energy density, isotropic scattering case

$$\mu \frac{dI}{d\tau} = -I + \frac{1}{2} \int_{-1}^{+1} I d\mu = -I + J$$

$$I = \int_0^{\infty} J(\tau') \exp\left(-\left|\frac{\tau' - \tau}{\mu}\right|\right) d\left(\left|\frac{\tau' - \tau}{\mu}\right|\right)$$

$$J(\tau) = \frac{1}{2} \int_0^{\infty} d\tau' J(\tau') E_1(|\tau - \tau'|)$$

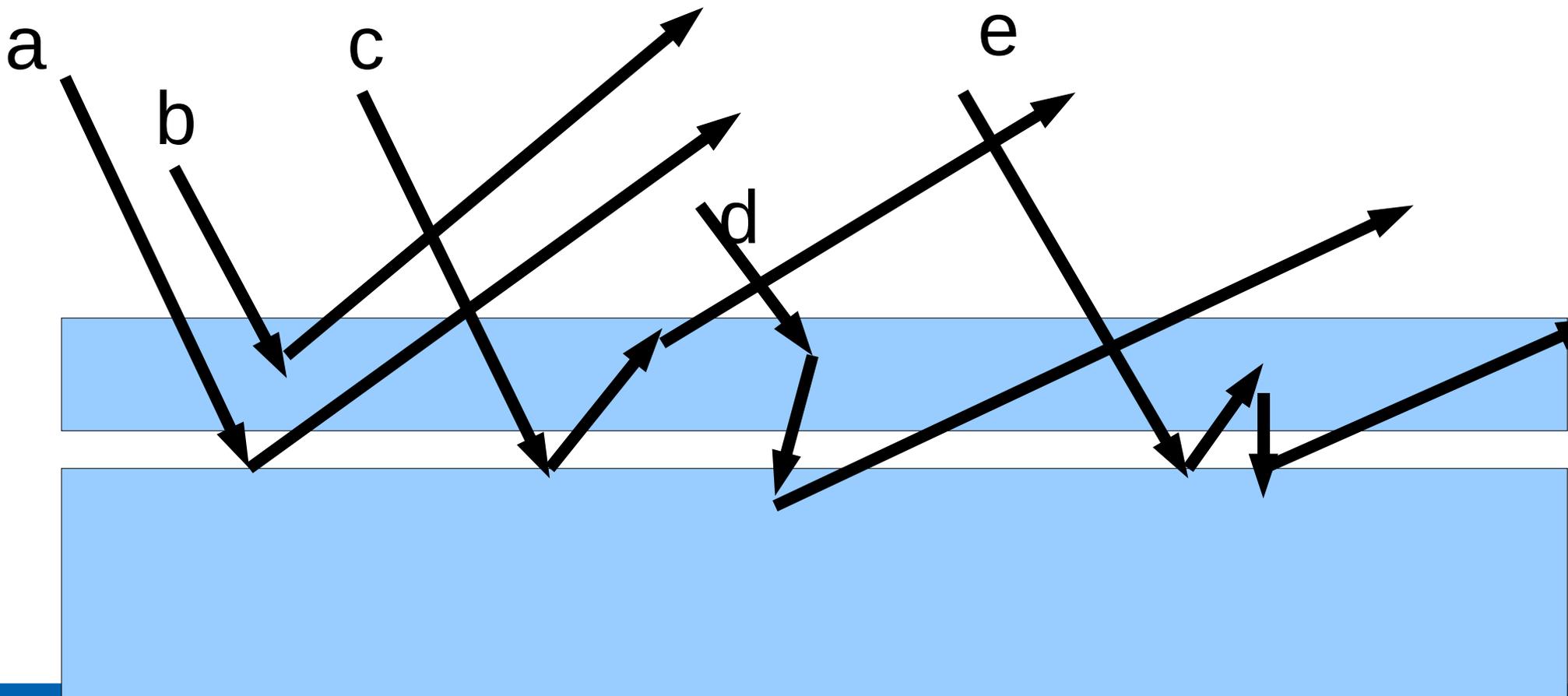
Becomes messier for anisotropic scattering!

The book ...

- *Radiative transfer* starts with 'discrete ordinates' i.e Gaussian quadrature (Wick) – seems a pedestrian exercise in coupled ODE's with constant coefficients until
- Exact and mysterious analytic results emerge from identities between zeros of Legendre polynomials and of the characteristic polynomial e.g $F = J/\sqrt{3}$; factorisation of the 'reflection probability'
$$S(\mu, \varphi, \mu_0, \varphi_0)$$
- They fall into place with the help of a physical principle of invariance (Ambartsumian)

Ambartsumian's principle of invariance, in pictures

All corrections due to adding a thin layer $d\tau$ must cancel !



....and in (not quite) full analytical
glory

$$\begin{aligned} S(\mu, \mu_0) \left(\frac{1}{\mu} + \frac{1}{\mu_0} \right) &= 1 + \frac{1}{2} \int S(\mu', \mu_0) d\mu' / \mu' \\ &+ \frac{1}{2} \int S(\mu, \mu'') d\mu'' / \mu'' \\ &+ \frac{1}{4} \int \int S(\mu, \mu'') S(\mu', \mu_0) \frac{d\mu'}{\mu'} \frac{d\mu''}{\mu''} \end{aligned}$$

Enter the H-function

$$H(\mu) = 1 + \frac{1}{2} \int S(\mu, \mu') d\mu' / \mu'$$

$$S(\mu, \mu_0) (1/\mu + 1/\mu_0) = H(\mu) H(\mu_0)$$

$$H(\mu) = 1 + \frac{\mu H(\mu)}{2} \int \frac{H(\mu')}{(\mu + \mu')} d\mu'$$

More complex problems like finite slabs, anisotropic scattering, polarisation all have their own H functions
Guiding principles enable C to navigate through these

The polarisation of skylight

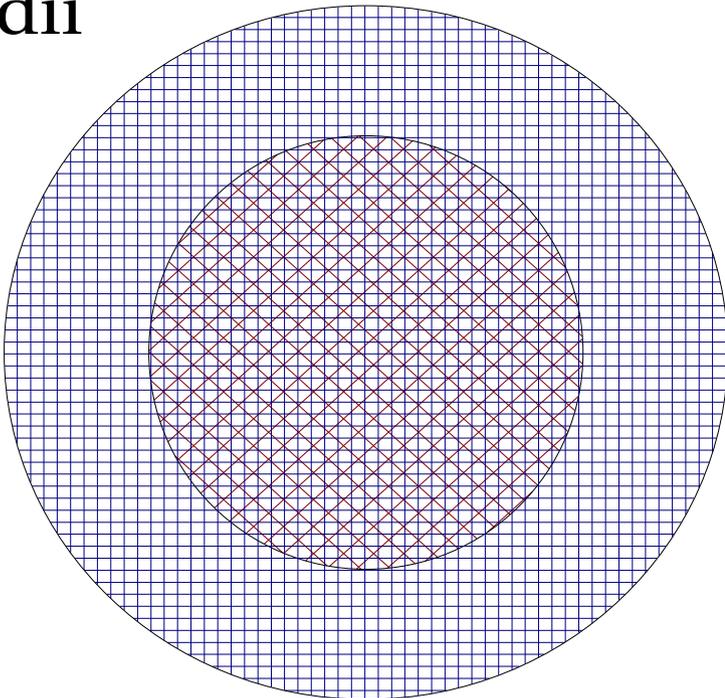
- Rayleigh's 1871 paper is single scattering, so maximum (100 per cent) polarisation at 90 degrees to the solar direction, zero at 0 degrees
- In reality, zero polarisation is displaced from the solar and antisolar directions, and maximum is about 90 per cent.
Qualitatively understood in early 20 century
- Chandrasekhar formulated a matrix transfer equation for the Stokes parameters, and found the 'eleven per cent solution' for stars and a quantitative fit to the sky

Chandrasekhar as quantum chemist

- Bethe, Hylleraas used a very complicated wave function to prove that the Hydride ion is bound
- Wildt realised the role of H^- in fitting the solar spectrum in near IR
- Chandrasekhar produced a much better energy from a much simpler wave function
- And followed it up with heroic numerical work on energies and oscillator strengths

The snowflake in hell: hydride ion in the Sun

Chandrasekhar's wave function for H⁻ reduces Coulomb repulsion by putting the two electrons in orbitals of different radii



Used different methods for computing transition probabilities from bound to continuum, and more elaborate wave functions

Some questions arising while reading RT

- Why replace one approximation by another?
- Why replace a linear integral equation by nonlinear ones?
- Why so much numerical work and tables?
- Why spend so much time and effort on a problem which was essentially solved?
- Why was it one of the 'happiest periods' of his scientific life?

(My) answers

- The loose ends left behind by Rayleigh scattering drew him to the field. Earlier methods and even formulations were inadequate (e.g Stokes parameters)
- The numerical work started out as such. His happy choice of Gaussian division and weights plus his rapport with equations drew him further and further into the problem
- The invariance principle deeply attracted him and he generalised it as far as he could
- His aim: not just answers, but a coherent structure, : RT syntactically closest to his ideal